

1 **North Atlantic Ocean-Originated Multicentennial Oscillation of the AMOC: A**  
2 **Coupled Model Study**

3  
4 Kunpeng Yang,<sup>a</sup> Haijun Yang,<sup>\*a</sup> Yang Li,<sup>b</sup> and Qiong Zhang<sup>c</sup>

5 <sup>a</sup> *Department of Atmospheric and Oceanic Sciences and Key Laboratory of Polar Atmosphere-ocean-*  
6 *ice System for Weather and Climate of Ministry of Education, Fudan University, Shanghai, 200438,*  
7 *China*

8 <sup>b</sup> *Department of Atmospheric and Oceanic Sciences, School of Physics, Peking University, Beijing,*  
9 *100871, China*

10 <sup>c</sup> *Department of Physical Geography and Bolin Centre for Climate Research, Stockholm University,*  
11 *Stockholm, 10691, Sweden*

12  
13 *Journal of Climate*

14 1<sup>st</sup> submission, July 13, 2023

15 1<sup>st</sup> revision, October 30, 2023

16  
17 *Corresponding author:* Haijun Yang, [yanghj@fudan.edu.cn](mailto:yanghj@fudan.edu.cn)

18

19  
20  
21  
22  
23  
24  
25  
26  
27  
28  
29  
30  
31  
32  
33  
34  
35  
36  
37  
38  
39

## ABSTRACT

A multicentennial oscillation (MCO) of the Atlantic meridional overturning circulation (AMOC) is exhibited in a CESM1 control simulation. It primarily arises from internal oceanic processes in the North Atlantic, potentially representing a North Atlantic Ocean-originated mode of AMOC multicentennial variability (MCV) in reality. Specifically, this AMOC MCO is mainly driven by salinity variation in the subpolar upper North Atlantic, which dominates local density variation. Salinity anomaly in the subpolar upper ocean is enhanced by the well-known positive salinity advection feedback that is realized through anomalous advection in the subtropical-subpolar upper ocean. Meanwhile, northward advection of salinity anomaly from the subtropical intermediate ocean through mean advection weakens the salinity anomaly in the subpolar upper ocean, leading to its phase change. This mechanism aligns with a theoretical model we proposed earlier. In this theoretical model, artificially deactivating either the anomalous or mean advection in the AMOC upper branch prevents it from exhibiting AMOC MCO, underscoring the indispensability of both the anomalous and mean advections in this North Atlantic Ocean-originated AMOC MCO. In our coupled model simulation, the South Atlantic and Southern Ocean do not exhibit variabilities synchronous with the AMOC MCO; the Arctic Ocean's contribution to the subpolar salinity anomaly is much weaker than the North Atlantic. Hence, this North Atlantic Ocean-originated AMOC MCO is distinct from the previously proposed Southern Ocean-originated and Arctic Ocean-originated AMOC MCOs.

**KEYWORDS:** Coupled model, Atlantic meridional overturning circulation, Multicentennial oscillation, Salinity anomaly

## 40 **1. Introduction**

41 Paleoclimatic proxy data from various sources reveal multicentennial variability (MCV) of the  
42 Earth's climate system (Chapman and Shackleton 2000; Nyberg et al. 2002; Kim et al. 2004; Wanner  
43 et al. 2008; Newby et al. 2014; Askjær et al. 2022), yet its origin and underlying mechanism remain  
44 elusive. The Atlantic meridional overturning circulation (AMOC), a crucial regulator of the climate  
45 system, also displays variability on this multicentennial timescale (Kissel et al. 2013; Thornalley et al.  
46 2013). Consequently, MCV of the AMOC has been proposed as a potential driver of the climate  
47 system's MCV (McDermott et al. 2001; Oppo et al. 2003; Hall et al. 2004; Miettinen et al. 2012;  
48 Chabaud et al. 2014; Ayache et al. 2018; Thirumalai et al. 2018).

49 Due to the scarcity of long-term direct AMOC observations, research on AMOC MCV has  
50 primarily relied on models of a hierarchy of complexities. Early studies have identified AMOC  
51 MCVs in two-dimensional (Mysak et al. 1993) and three-dimensional (Mikolajewicz and Maier-  
52 Reimer 1990; Winton and Sarachik 1993; Drijfhout et al. 1996) ocean models. Over the past two  
53 decades, several studies have utilized the more sophisticated coupled models to study AMOC MCV  
54 (Park and Latif 2008; Delworth and Zeng 2012; Martin et al. 2013, 2015; Jiang et al. 2021; Meccia et  
55 al. 2023). Perhaps because of the differences in models employed, the mechanisms underlying these  
56 coupled model AMOC MCVs vary. Consequently, these studies might have identified distinct  
57 potential modes of the real-world AMOC MCV, which do not necessarily contradict each other.  
58 Hence, it might be more suitable to refer to the MCVs of AMOC in these studies as MCOs, as the  
59 term "variability" is usually from an observational or statistical view, but "oscillation" is often  
60 dominated by a specific dynamic mode and carries a more physical connotation (Liu 2012; Sutton et  
61 al. 2018; Zhang et al. 2019). In this context, studying the AMOC MCV essentially involves tackling a  
62 dynamic system problem, that is, a comprehensive understanding of the real-world AMOC MCV  
63 critically depends on understanding each constituent mode (or AMOC MCO). Analyses of the  
64 modeled AMOC MCOs, therefore, should draw insights from the more fundamental (and typically  
65 earlier) theoretical studies. However, these coupled model studies often do not interpret their AMOC  
66 MCOs from a more theoretical and dynamic perspective, highlighting a disconnect between model  
67 results and theories.

68 Unlike the El Niño-Southern Oscillation, where both linear (Suarez and Schopf 1988; Jin 1997)  
69 and nonlinear theories (Tziperman et al. 1994; Sun 1997) have been extensively developed, the  
70 majority of low-frequency AMOC oscillation theories focus on linear oscillation (Griffies and

71 Tziperman 1995; Rivin and Tziperman 1997; Wei and Zhang 2022), where the AMOC oscillation is  
72 regular and symmetric. Under the linear framework, the single-equilibrium oscillation is perceived as  
73 an anomaly hovering around the unstable equilibrium. Positive and negative feedbacks enhance and  
74 weaken the anomaly, collectively leading to the anomaly's phase transition and therefore its cyclic  
75 evolution. This forms our foundational "theoretical interpretation" of linear oscillation, by which we  
76 are inspired to review the aforementioned coupled model studies on AMOC MCO.

77 In an earlier study using the Kiel Climate Model (KCM), Park and Latif (2008) found an AMOC  
78 oscillation with a period of 300-400 years. Their follow-up studies, Martin et al. (2013, 2015)  
79 proposed that this AMOC MCO originates from the Southern Ocean. When the AMOC is  
80 anomalously strong, heat content of the mid-depth water in the Weddell Sea increases due to  
81 strengthened southward transport of the warmer North Atlantic Deep Water (NADW). Deep  
82 convection in the Southern Ocean is triggered when the mid-depth heat accumulation becomes  
83 excessive, hence the warm deeper water reaches the cold surface air and convectively releases heat to  
84 the atmosphere, densifying the Weddell Sea overall. Therefore, the Atlantic north-to-south density  
85 gradient is decreased, limiting the NADW formation and thus the AMOC strength (Hughes and  
86 Weaver 1994). This is similar to the advective-convective mechanism proposed by Yin (1995), where  
87 convection is initiated by the advective buildup of heat. Therefore, the oscillation timescale is set by  
88 the advective heat accumulation. Yet, this AMOC MCO is induced by the drastic "flip-flop"  
89 convection (Welander 1982) in the Southern Ocean, which is in essence a multi-equilibrium  
90 phenomenon.

91 Using a GFDL CM2.1 model simulation, Delworth and Zeng (2012) also identified an AMOC  
92 MCO related to the Southern Ocean, but with a different mechanism. Starting with a weak AMOC, a  
93 positive surface salinity anomaly in the Southern Ocean is carried northward by the mean circulations  
94 in the upper branch of the AMOC, strengthening the NADW formation when it reaches the North  
95 Atlantic convection region, driving the AMOC into its positive phase. Synchronously, more  
96 freshwater is produced in the Southern Ocean due to the positive AMOC anomaly, and would be  
97 transported northward later. This will weaken the AMOC when the negative salinity anomaly reaches  
98 the northern convection region, completing a full cycle. The oscillation timescale here is determined  
99 by the time consumed in transporting the Southern Ocean salinity anomaly to the North Atlantic; no  
100 drastic variation of deep convection in the Southern Ocean is documented by the authors. Therefore,  
101 it is a Southern Ocean-originated AMOC MCO whose essence differs from the "flip-flop" AMOC

102 MCO in Park and Latif (2008). In short, the mean advection process is raised as the driver for the  
103 entire cycle, including the growing and weakening of the anomalies, as well as the phase transition.  
104 However, considering that the mean advection process is a weakening process for salinity anomaly in  
105 the subpolar North Atlantic (Griffies and Tziperman 1995; Wei and Zhang 2022), there should be  
106 processes that enhance the subpolar salinity anomaly, which are not resolved in this study.

107 Recently, a group of studies identified Arctic Ocean-originated AMOC MCOs. Jiang et al. (2021)  
108 found a 200-year AMOC MCO in their IPSL-CM6A-LR model simulation. When the AMOC resides  
109 in its strong phase, the Arctic Ocean is warmed and thus more sea ice melting leads to negative  
110 salinity anomaly therein. The negative Arctic Ocean salinity anomaly is advected southward through  
111 mean advection, inhibiting the subpolar deep convection and driving the AMOC into its negative  
112 phase. Meccia et al. (2023) found a 150-year AMOC oscillation in the EC-Earth3 model with a  
113 similar mechanism. Resemblance in explanations proposed by these two studies might be attributed to  
114 the shared ocean component (NEMO 3.6) of their models. As with Delworth and Zeng (2012), the  
115 mean advection is again proposed as the process governing the entire evolution, suggesting that  
116 enhancing processes for the AMOC anomaly remain to be found. Another study by Mehling et al.  
117 (2023) also proposed salinity anomaly from the Arctic Ocean as the driver for their modeled AMOC  
118 MCO, but they utilized an intermediate-complexity model instead of a high-complexity coupled  
119 model. In addition, Vellinga and Wu (2004) analyzed an AMOC oscillation on centennial instead of  
120 multicentennial timescale, with the air-sea interaction rather than internal oceanic processes as the  
121 core mechanism. As such, we have not included Vellinga and Wu (2004) and Mehling et al. (2023) in  
122 coupled model AMOC MCO studies.

123 The ability to test whether these aforementioned AMOC MCOs exist in reality is constrained by  
124 observational limitations. Therefore, interpreting coupled model AMOC MCOs from a more  
125 fundamental and theoretical perspective is a practical approach to improve the understanding of the  
126 real-world AMOC MCV at this stage. By analyzing results from a CESM1 control simulation, we  
127 identified an AMOC MCO dominated mainly by processes in the North Atlantic. Its mechanism can  
128 be explained by a linear AMOC MCO theory we proposed earlier in Li and Yang (2022) (hereafter  
129 LY22). The theoretical box model in LY22 is inspired by another theoretical box model proposed by  
130 Griffies and Tziperman (1995) (hereafter GT95), which focuses on AMOC multidecadal oscillation.  
131 Both theoretical models only include processes in the North Atlantic, and are therefore unable to  
132 exhibit AMOC oscillations related to other ocean basins. In the GT95 theoretical model, the depth of

133 the upper boxes is set at only 300 m. In addition, the volume of the subpolar boxes is set to only 1/11  
134 of the North Atlantic, which is too small as the latitude for deep water formation (DWF) spans from  
135 approximately 50°N to 75°N (Hansen et al. 2003; Kieke and Yashayaev 2015). LY22 adopted new  
136 model parameters including thicker upper boxes and larger subpolar boxes, and found that the model  
137 exhibits AMOC MCO. By including a nonlinear subpolar vertical mixing into the purely linear  
138 advection system, a self-sustained AMOC oscillation is realized in LY22. By contrast, GT95 focused  
139 on stochastically-forced oscillation and the AMOC multidecadal oscillation therein is a damped  
140 oscillation. In LY22, the AMOC anomaly is enhanced by the positive salinity advection feedback in  
141 the upper level of the subtropical-subpolar North Atlantic, and weakened by the mean advection of  
142 salinity anomaly in the upper ocean. Starting with a positive AMOC anomaly, more subtropical saline  
143 water is transported northward, further strengthening the AMOC anomaly. This constitutes the  
144 classical positive salinity advection feedback (Stommel 1961; Marotzke 1996). On the other hand, a  
145 negative salinity anomaly formed in the subtropical upper ocean is transported northward through  
146 mean advection, weakening the AMOC anomaly.

147 In this study, we will treat the AMOC MCO in the CESM1 control simulation as a linear  
148 oscillation, and analyze its mechanism in a linear framework. In section 2, an approach for extracting  
149 low-frequency variability from the raw model data is introduced. In section 3, the MCOs of the  
150 AMOC and global buoyancy fields are presented; dominance of salinity variation in the AMOC MCO  
151 is highlighted. In section 4, the evolution patterns of salinity anomalies are shown. In section 5,  
152 processes contributing to the AMOC MCO are quantitatively analyzed, and the significance of the  
153 key advection processes is examined employing the theoretical model in LY22. In section 6, the  
154 conclusion and discussion are provided.

155

## 156 **2. Model and methods**

### 157 *a. Pre-industrial control simulation*

158 The coupled model employed is the Community Earth System Model (CESM, version 1.0.4)  
159 developed by the National Centre for Atmospheric Research (NCAR). It is a global climate model  
160 consisting of five components: atmosphere, land surface, ocean, sea ice, and ice sheet (not active). A  
161 coupler exchanges data between these components.

162 The model grid utilized is f19\_gx1v6. The atmospheric component is the Community Atmosphere  
163 Model version 4 (CAM4) (Neale et al. 2010), with 26 vertical levels and a horizontal resolution of  
164  $1.9^\circ \times 2.5^\circ$ . The land surface model is the Community Land Model version 4 (CLM4) (Lawrence et  
165 al. 2011), with the same horizontal resolution as CAM4. The ocean model is the Parallel Ocean  
166 Program version 2 (POP2) (Smith et al. 2010). It uses the gx1v6 curvilinear grid, having  $384 \times 320$   
167 grid points horizontally and 60 layers vertically. The horizontal grid is zonally uniform at a  $1.125^\circ$   
168 resolution but meridionally non-uniform, with a  $0.27^\circ$  resolution near the equator, increasing to  $0.65^\circ$   
169 at  $60^\circ\text{N/S}$  and then decreasing toward the polar regions. The sea ice model is the Community Ice  
170 Code (CICE4) (Hunke and Lipscomb 2010), with the same horizontal resolution as POP2. The  
171 coupler is the CESM Coupler CPL7 (Craig et al. 2012).

172 In the ocean model, velocity is divided into three components: explicit Eulerian-mean velocity,  
173 parameterized bolus velocity, and parameterized sub-mesoscale velocity (Gent and McWilliams 1990;  
174 Fox-Kemper and Ferrari 2008; Fox-Kemper et al. 2008); the latter two are collectively regarded as the  
175 parameterized eddy-induced velocity. These three velocity components have their corresponding  
176 transport or streamfunction. The total, or referred to as “residual,” velocity (streamfunction) is the  
177 sum of these three components. A 2500-year control simulation is conducted from the rest with pre-  
178 industrial configuration, to assure that the thermohaline circulation has reached its equilibrium before  
179 our study period. In this study, outputs of the last 1500 years are analyzed.

180

## 181 *b. Data analysis methods*

182 Instead of using the traditional empirical orthogonal function (EOF) method, the low-frequency  
183 component analysis (LFCA) method (Wills et al. 2018) is adopted to evaluate the low-frequency  
184 AMOC variability. The LFCA provides linear combinations of the first  $n$  EOFs and principal  
185 components (PCs) of the data. The results are  $n$  low-frequency patterns (LFPs) as the spatial patterns  
186 and  $n$  corresponding low-frequency components (LFCs) as the time series. The LFPs and  
187 corresponding LFCs are ranked in a descending order according to the ratio of their low-frequency  
188 variances (obtained through a low-pass filter) to the total variance of the first  $n$  EOFs. Therefore, low-  
189 frequency signals are the most concentrated in LFP1 and LFC1, which are our focus. More details  
190 about the LFCA method can be found in Wills et al. (2018). As demonstrated in Jiang et al. (2021),

191 the LFCA outperforms the EOF method in extracting and analyzing low-frequency variability of the  
192 AMOC.

193 A Lanczos filter with 121 weights and 60 years cut-off period is used for the LFCA and for  
194 filtering other time series throughout this study. Power spectral analysis is conducted for both the  
195 unfiltered AMOC index and AMOC LFC1. Significance at 95% confidence level and the best-fit  
196 first-order Markov red noise spectrum are presented along with the power spectrum.

197 To quantify the contributions to salinity anomaly of a specific region from different processes,  
198 liquid freshwater mass transport ( $\text{kg s}^{-1}$ ) into the region is computed. The total liquid freshwater mass  
199 transports across the x-direction and y-direction of the ocean model grids are computed as:

$$200 \quad F_x = \int_{y_s}^{y_n} \int_{z_b}^{z_t} \left(1 - \frac{S}{S_{ref}}\right) \rho u dz dy \quad (1)$$

$$201 \quad F_y = \int_{x_w}^{x_e} \int_{z_b}^{z_t} \left(1 - \frac{S}{S_{ref}}\right) \rho v dz dx \quad (2)$$

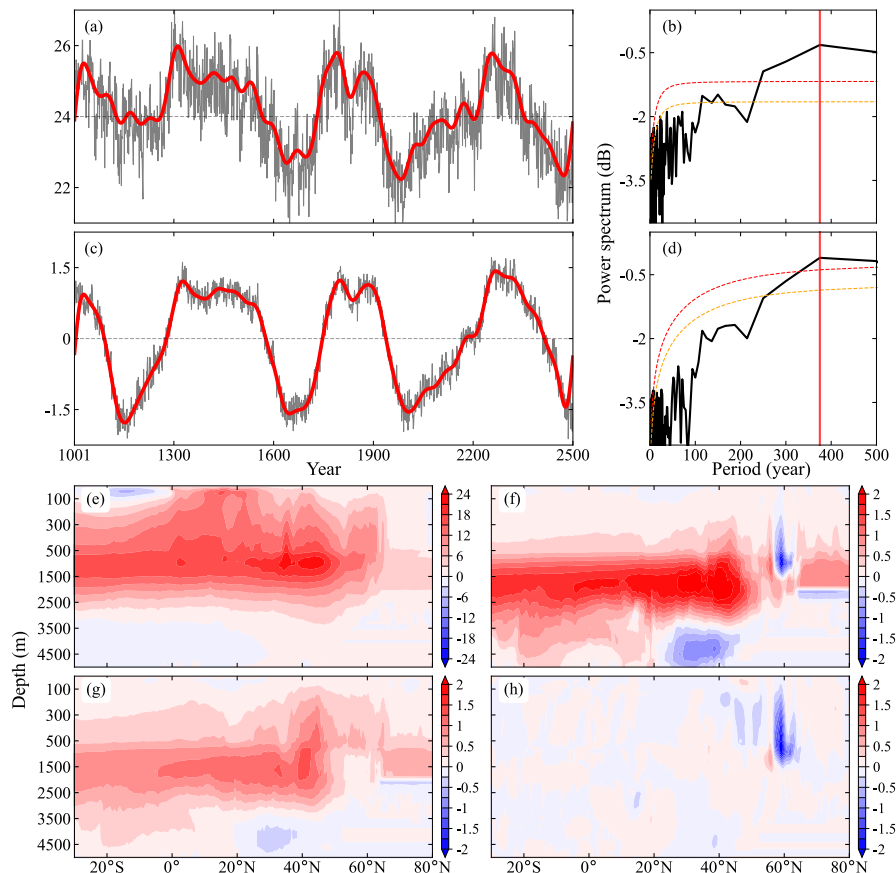
202 , respectively. The anomaly of advection-induced liquid freshwater mass transport can be linearly  
203 decomposed into three components: (1) anomalous advection of mean salinity induced; (2) mean  
204 advection of salinity anomaly induced; and (3) anomalous advection of salinity anomaly induced. The  
205 third component is a nonlinear advection term that is much smaller than the former two components,  
206 hence is disregarded. The liquid freshwater mass transports induced by anomalous advection of mean  
207 salinity are expressed as  $\int_{y_s}^{y_n} \int_{z_b}^{z_t} \left(1 - \frac{\bar{S}}{S_{ref}}\right) \rho u' dz dy$  and  $\int_{x_w}^{x_e} \int_{z_b}^{z_t} \left(1 - \frac{\bar{S}}{S_{ref}}\right) \rho v' dz dx$ ; the liquid  
208 freshwater mass transports induced by mean advection of salinity anomaly are expressed as  
209  $\int_{y_s}^{y_n} \int_{z_b}^{z_t} -\frac{S'}{S_{ref}} \rho \bar{u} dz dy$  and  $\int_{x_w}^{x_e} \int_{z_b}^{z_t} -\frac{S'}{S_{ref}} \rho \bar{v} dz dx$ .  $y_s$ ,  $y_n$ ,  $x_w$ , and  $x_e$  are the southernmost,  
210 northernmost, westernmost, and easternmost grid points.  $z_b$  and  $z_t$  are the bottommost and uppermost  
211 grid points.  $S$ ,  $\bar{S}$ , and  $S'$  are the three-dimensional sea water salinity in psu, its climatological value,  
212 and its anomaly.  $S_{ref}$  is the reference sea water salinity of the given study area.  $\rho$  is the sea water  
213 density in  $\text{kg m}^{-3}$ .  $u$ ,  $\bar{u}$ , and  $u'$  ( $v$ ,  $\bar{v}$ , and  $v'$ ) are the three-dimensional x-direction (y-direction)  
214 velocity in  $\text{m s}^{-1}$ , its climatological value, and its anomaly.



215

216 **3. MCOs in the coupled model**217 *a. AMOC*

218 The modeled AMOC index exhibits a distinct MCV (Fig. 1a), with the most significant peak  
 219 around 375 years (Fig. 1b). The AMOC index is defined as the maximum total meridional  
 220 streamfunction within the North Atlantic region spanning 20°-70°N and 200-3000 m. The AMOC  
 221 index exhibits a stable oscillation around its mean state, with a magnitude of around 2 Sv, about 10%  
 222 of the climatological value (24 Sv). As we will interpret the modeled AMOC MCV as a linear  
 223 oscillation, henceforth the MCVs of the AMOC and other variables will be referred to as MCOs. The  
 224 climatological AMOC exhibits an overall northward branch in the upper 0-1000 m, a deep convection  
 225 branch around 60°N, and a southward NADW branch in the deep ocean of 1500-3000 m (Fig. 1e).  
 226 The maximum value is located near 1000 m at around 40°N.



227

228 FIG. 1. (a) Time series for the Atlantic meridional overturning circulation (AMOC) index (units: Sv;  $1 \text{ Sv} = 10^6$   
229  $\text{m}^3 \text{ s}^{-1}$ ) of model years 1001-2500. The AMOC index is defined as the maximum total meridional streamfunction in  
230 the North Atlantic spanning  $20^\circ$ - $70^\circ\text{N}$  and 200-3000 m. The gray curve represents the unfiltered AMOC index, and  
231 the red curve is the low-pass-filtered AMOC index using the Lanczos filter. The horizontal dashed line denotes the  
232 climatological value of the AMOC (24 Sv). (b) Power spectrum (units: dB) of the unfiltered AMOC index, with  
233 period as the abscissa. The dashed orange and red curves represent the best-fit first-order Markov red noise spectrum  
234 and the significance at 95% confidence level, respectively. The vertical red line denotes the most significant peak  
235 (375 years). (c) Same as (a), but for the AMOC's first low-frequency component (LFC1). Before applying the low-  
236 frequency component analysis (LFCA) method, the data is detrended and then weighted according to the square root  
237 of grid cell thicknesses. The Lanczos filter is used for the LFCA. (d) Same as (b), but for the AMOC LFC1. (e)  
238 Climatological pattern of the total AMOC averaged over years 1001-2500 (units: Sv). (f) Pattern of the AMOC's  
239 first low-frequency pattern (LFP1) (units: Sv). (g) and (h) are the regression patterns of the Eulerian-mean and eddy-  
240 induced AMOCs on the AMOC LFC1 (units: Sv), respectively.

241

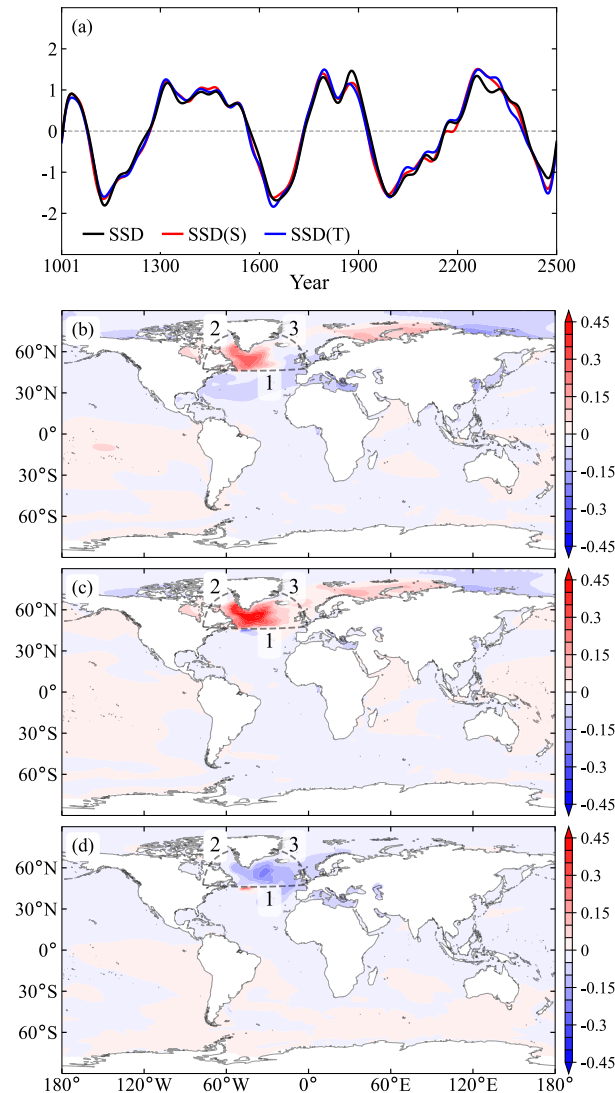
242 To better depict the low-frequency variability of the AMOC, the LFCA method is applied to the  
243 total AMOC. The first 10 EOFs are utilized, which explain 93.7% of the total variance. High-  
244 frequency signals of the AMOC are substantially weakened in the AMOC LFC1 (Figs. 1c, d). The  
245 power spectrum of the AMOC LFC1 reveals that the most significant peak persists near 375 years and  
246 becomes the only peak passing the significance test (Fig. 1d). The AMOC LFP1 accounts for 87.5%  
247 of the low-frequency variance of the first 10 EOFs. The pattern of the AMOC LFP1 (Fig. 1f) aligns  
248 with that of the climatological AMOC (Fig. 1e), but exhibits a structure with negligible transport  
249 above 500 m. The upper ocean wind-driven subtropical cell, which is discernible in the climatological  
250 AMOC pattern (Fig. 1e), has been filtered out in the AMOC LFP1 (Fig. 1f). This suggests that the  
251 AMOC MCO in this study primarily occurs in the lower ocean, and the high-frequency variability  
252 that may be related to the upper ocean wind-driven circulation should be ruled out.

253 The Eulerian-mean component of the AMOC demonstrates a coherent meridional variability  
254 throughout the Atlantic basin (Fig. 1g), while the eddy-induced component has a significant local  
255 variability within the subpolar North Atlantic (Fig. 1h). The fluctuation of the eddy-induced AMOC is  
256 of comparable magnitude to that of the Eulerian-mean AMOC, but with an opposite sign in the  
257 subpolar deep convection region, resulting in the negative signal at the same position of the AMOC  
258 LFP1 (Fig. 1f). This suggests that a stronger (weaker) Eulerian-mean AMOC is associated with a  
259 weaker (stronger) eddy-induced AMOC. Although the causality between the Eulerian-mean and  
260 eddy-induced AMOCs in the current study is not yet clear, Figs. 1f and 1h suggest that the eddy-  
261 induced AMOC plays a role in this AMOC MCO.

262

263 *b. Global surface buoyancy fields*

264 Consistent with the AMOC LFC1, the global ocean's surface buoyancy fields also exhibit  
 265 remarkable MCOs (Fig. 2). The LFCA is conducted for the global sea surface density (SSD) anomaly,  
 266 SSD anomaly induced by sea surface salinity (SSS) anomaly, and SSD anomaly induced by sea  
 267 surface temperature (SST) anomaly according to Roquet et al. (2015). All of their LFC1s reflect  
 268 variability on multicentennial timescale (Fig. 2a).



269

270 FIG. 2. (a) Filtered LFC1s (units: dimensionless) of the global sea surface density (SSD) anomaly (black curve),  
 271 SSD anomaly induced by sea surface salinity (SSS) anomaly (red curve), and SSD anomaly induced by sea  
 272 surface temperature (SST) anomaly (blue curve). (b) LFP1 (units:  $\text{kg m}^{-3}$ ) of the global SSD anomaly. (c) and (d) are  
 273 the same as (b), but for the SSD anomalies induced by SSS and SST anomalies, respectively. Before the LFCA, the data  
 274 is detrended and then weighted according to the square root of grid cell areas. The Lanczos filter is applied in (a);  
 275 and the LFCA, in (b)-(d). The region enclosed by boundaries 1-3 (dashed curves in Figs. 2b-d) represents the deep  
 276 water formation (DWF) region, and will be used later. These boundaries are parallel to the grid lines of the ocean  
 277 model. Boundary 1 is along 47°N.

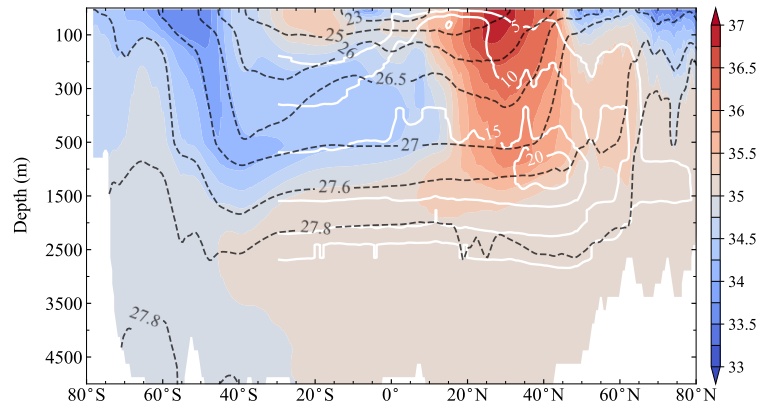
278

279 The LFP1s of the global surface buoyancy fields indicate that the strongest multicentennial signals  
280 are located in the North Atlantic (Figs. 2b-d), especially in the DWF region in the subpolar North  
281 Atlantic with the deepest March mixed layer depth simulated (figure not shown). These spatial  
282 patterns indicate that in both the North Atlantic and Arctic Ocean, the SSD anomalies (Fig. 2b) are  
283 dominated by SSS anomalies (Fig. 2c), yet partly canceled by SST anomalies (Fig. 2d). Over the  
284 “transition zone” near 45°N along the Gulf Stream extension (Buckley and Marshall 2016), the local  
285 density variability is negligible (Fig. 2b) due to the counteractive effects from anomalous salinity and  
286 temperature. Multicentennial signals in the other basins, such as the South Atlantic, Southern Ocean,  
287 and the Pacific and Indian Oceans, are rather weak. In this study, the DWF region is defined as the  
288 region enclosed by boundaries 1-3 (Figs. 2b-d); boundary 1 is situated just to the “transition zone.”  
289

## 290 **4. Evolution of salinity anomalies in the Atlantic**

### 291 *a. Latitude-depth patterns*

292 We first examine the model distribution of climatological salinity in the Atlantic (Fig. 3). In the  
293 North Atlantic, the meridional gradient of climatological salinity is generally greater in the upper  
294 ocean and decreases with depth. Specifically, the most saline water (salinity higher than 36.5 psu) is  
295 located in the subtropical upper ocean between 20°N and 40°N, extending downward to 1500 m and  
296 then southward to 40°S in the deep ocean. The more saline subtropical water is separated from the  
297 fresher subpolar water by the Gulf Stream extension and the North Atlantic Current (NAC), forming  
298 the subpolar front near 45°N, whose location is consistent with the downward branch of the AMOC  
299 (Fig. 1e). The more saline subtropical water also extends northward to the subpolar subsurface ocean,  
300 against the freshwater from the Arctic Ocean. On the other hand, fresher water (salinity lower than 34  
301 psu) comes mainly from the surface of the subpolar Southern Ocean, extending downward and  
302 northward and occupying the southern subtropical ocean between 200 m and 1500 m, forming the  
303 Antarctic Intermediate water.



304

305 FIG. 3. Climatological salinity zonally averaged in the Atlantic (shading; units: psu), superimposed with  
 306 climatological potential density ( $\sigma_\theta$ ) (black contour; units:  $\text{kg m}^{-3}$ ) and AMOC (white contour; units: Sv).

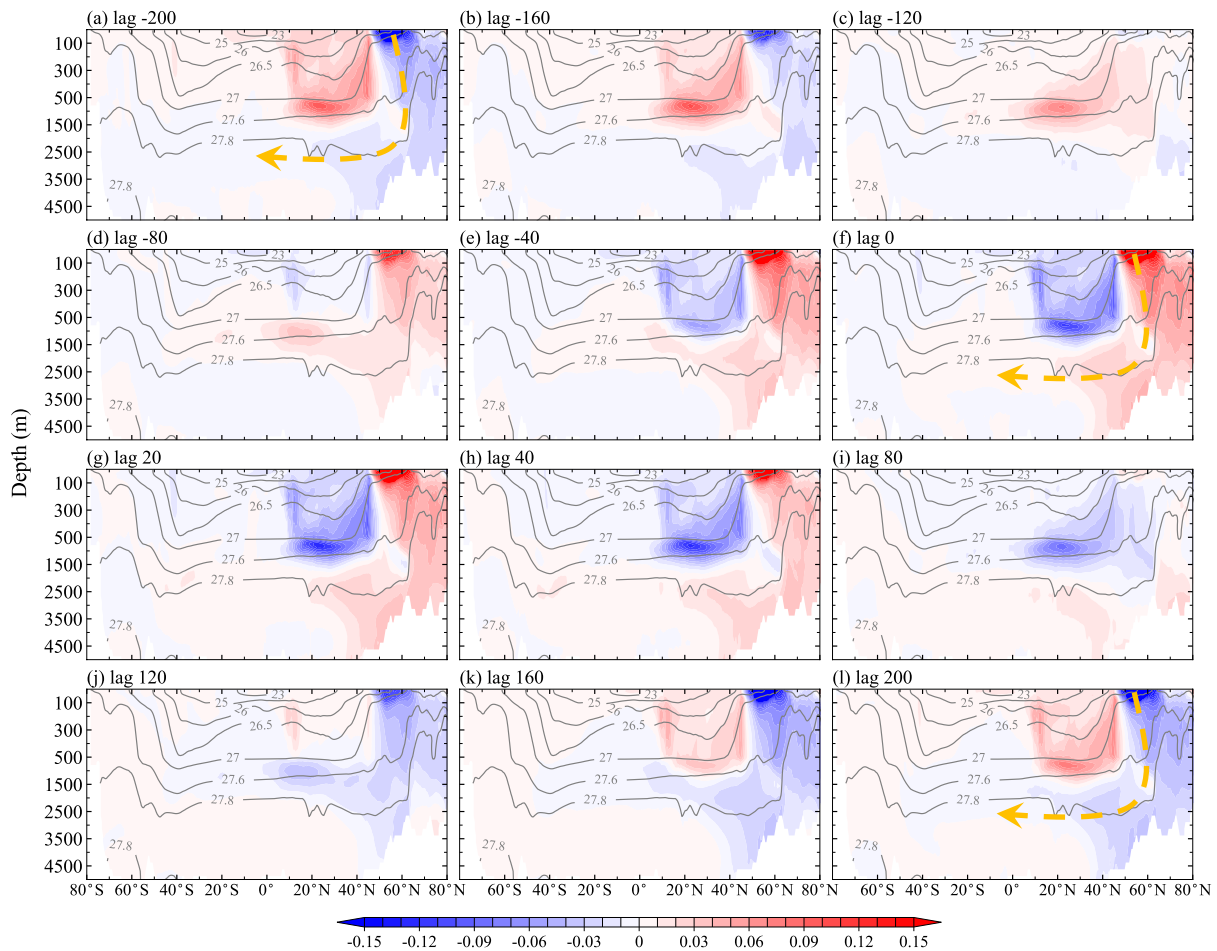
307

308 Next, evolution of salinity anomalies with the AMOC over the multicentennial cycle is analyzed.  
 309 Figure 4 illustrates the lead/lag regression coefficients of the zonally averaged salinity anomalies in  
 310 the Atlantic on the AMOC LFC1, superimposed with climatological potential density ( $\sigma_\theta$ ). Positive  
 311 and negative regression coefficients at lag  $n$  years represent that generally there are positive and  
 312 negative salinity anomalies in the corresponding regions, respectively, when salinity anomalies lag  
 313 the AMOC LFC1 by  $n$  years. For conciseness, positive/negative salinity anomaly is used to represent  
 314 positive/negative regression coefficient, which can be rough to some extent.

315 When the salinity anomalies lead the AMOC LFC1 by 200 years (Fig. 4a), there is a pronounced  
 316 negative salinity anomaly centered in the upper ocean around  $55^\circ\text{N}$ , corresponding to the weakest  
 317 NADW formation and AMOC. This negative anomaly extends from the surface to deep ocean of the  
 318 North Atlantic north of  $45^\circ\text{N}$ , and occupies 1500–3500 m in the deep ocean south of  $45^\circ\text{N}$ . South of  
 319 the subpolar negative anomalies, broad positive anomalies are observed at lower latitudes, occupying  
 320 the upper 1500 m of the Atlantic. The negative anomalies are the strongest in the upper DWF region,  
 321 while the positive anomalies have the greatest magnitude in the subtropical intermediate ocean  
 322 between 500 m and 1500 m. This dipole structure is the most robust feature throughout the entire  
 323 evolution of salinity anomalies in the North Atlantic.

324 The evolution of salinity anomalies at the multicentennial timescale is closely linked to the  
 325 AMOC's evolution. The downward and southward movements of salinity anomalies north of  $45^\circ\text{N}$   
 326 correspond to a strong convection or vertical mixing, and the mean advection by the lower branch of  
 327 the AMOC, respectively. The northward and upward movements of anomalies south of  $45^\circ\text{N}$  go

328 roughly within  $26.5\text{-}27.6 \sigma_\theta$ , corresponding to the mean advection through the upper branch of the  
 329 AMOC. These two anomalies circulate in the North Atlantic, changing their phases during their  
 330 movements (Fig. 4). In the South Atlantic, salinity anomalies do not reflect evolution synchronous  
 331 with the AMOC.



332

333 FIG. 4. Lead/lag regression coefficients of zonally averaged salinity anomalies in the Atlantic on the AMOC  
 334 LFC1 (shading; units: psu). Negative lag means the AMOC LFC1 lags the salinity anomalies (units: year). Contours  
 335 show the zonally averaged climatological potential density  $\sigma_\theta$  in the Atlantic (units:  $\text{kg m}^{-3}$ ). Orange arrows in (a),  
 336 (f), and (l) show schematically the downward and southward movements of salinity anomalies.

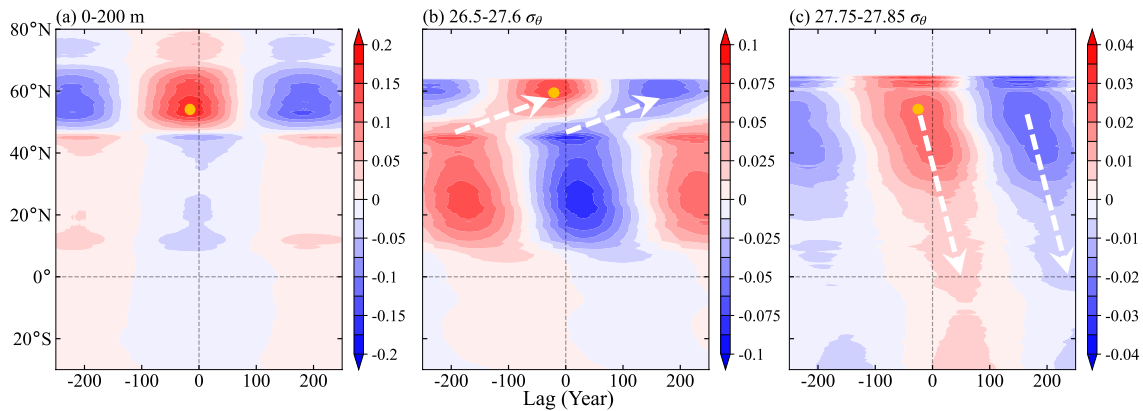
337

338 Specifically, during the period when the salinity anomalies lead the AMOC by 200-120 years  
 339 (Figs. 4a-c), salinity anomalies north of  $45^\circ\text{N}$  are negative and the AMOC is in its weak regime. The  
 340 magnitude of negative anomalies decreases with time, indicating an ongoing phase transition from the  
 341 weak to the strong AMOC regime. The positive anomaly in the subtropical intermediate ocean moves  
 342 northward from about  $40^\circ\text{N}$  and upward along  $26.5\text{-}27.6 \sigma_\theta$  (Fig. 4c), neutralizing the negative  
 343 anomaly in the DWF region and furthering its phase shift to positive (Fig. 4d). When the salinity

344 anomalies lead the AMOC by 80-0 years, positive anomalies north of 45°N develop gradually (Figs.  
345 4d-f) and eventually reach the maximum magnitude at lag 0 years (Fig. 4f), exhibiting a similar  
346 evolution to that of the AMOC. Anomaly in the upper DWF region is transported downward through  
347 convection or vertical mixing, then propagates southward in the deep ocean (Fig. 4, orange arrows).  
348 Meanwhile, negative anomalies grow in lower latitudes at 26.5-27.6  $\sigma_\theta$ . Afterward, the AMOC starts  
349 to decrease and the evolution of salinity anomalies enters the opposite phase (Figs. 4g-l). Throughout  
350 the entire cycle, salinity anomalies in the Arctic Ocean and Atlantic deep ocean are largely  
351 synchronized with that in the upper DWF region. Signals in the South Atlantic are much weaker  
352 compared to those in the North Atlantic.

353 Overall, the major evolution characteristics of salinity anomalies in the North Atlantic are  
354 reflected in three levels: (i) the upper ocean around 0-200 m, where the DWF region salinity anomaly  
355 is the strongest and develops locally, (ii) the intermediate ocean around 26.5-27.6  $\sigma_\theta$ , where the  
356 salinity anomalies south of 45°N evolve, and propagate northward to weaken the salinity anomaly in  
357 the upper DWF region, and (iii) the deep ocean around 27.8  $\sigma_\theta$ , where the salinity anomalies originate  
358 from the upper DWF region and propagate southward.

359 Figure 5 clearly shows the local development of salinity anomaly in the upper DWF region and  
360 the northward (southward) propagation of salinity anomaly in the intermediate (deep) ocean. Based  
361 on Fig. 4, we vertically average the anomalies over these three levels and then calculate their lead/lag  
362 regression coefficients on the AMOC LFC1. In the upper ocean (Fig. 5a), the most remarkable signal  
363 is in 45°-65°N, showing a local periodic evolution without a robust connection with signals in both  
364 the subtropical and polar regions. In the 26.5-27.6  $\sigma_\theta$  intermediate ocean (Fig. 5b), anomaly at 45°N  
365 propagates northward (white arrows), suggesting its potential influence on the DWF region. Anomaly  
366 near 20°N appears to develop locally with the opposite sign to that north of 45°N. In the 27.75-27.85  
367  $\sigma_\theta$  deep ocean (Fig. 5c), anomalies propagate southward from the subpolar to equatorial and South  
368 Atlantic (white arrows). In all these three levels, the maximum regression coefficient in the subpolar  
369 region occurs when the salinity anomalies lead the AMOC LFC1 by around 10 years (Fig. 5, orange  
370 dot).



371

372 FIG. 5. Lead/lag regression coefficients of zonally and vertically averaged salinity anomalies in the Atlantic on  
 373 the AMOC LFC1 (units: psu). (a) Averaged over 0-200 m, (b) averaged over 26.5-27.6  $\sigma_\theta$ , and (c) averaged over  
 374 27.75-27.85  $\sigma_\theta$ . The orange dot denotes the position of the maximum regression coefficient. In (b) and (c), salinity  
 375 anomalies shallower than 200 m or north of 65°N are removed. The dashed white arrows show schematically the  
 376 meridional propagations of salinity anomalies in the intermediate-deep oceans. Note that the colorbars for the three  
 377 subplots are different. Negative lag means the AMOC LFC1 lags the salinity anomalies (units: year).

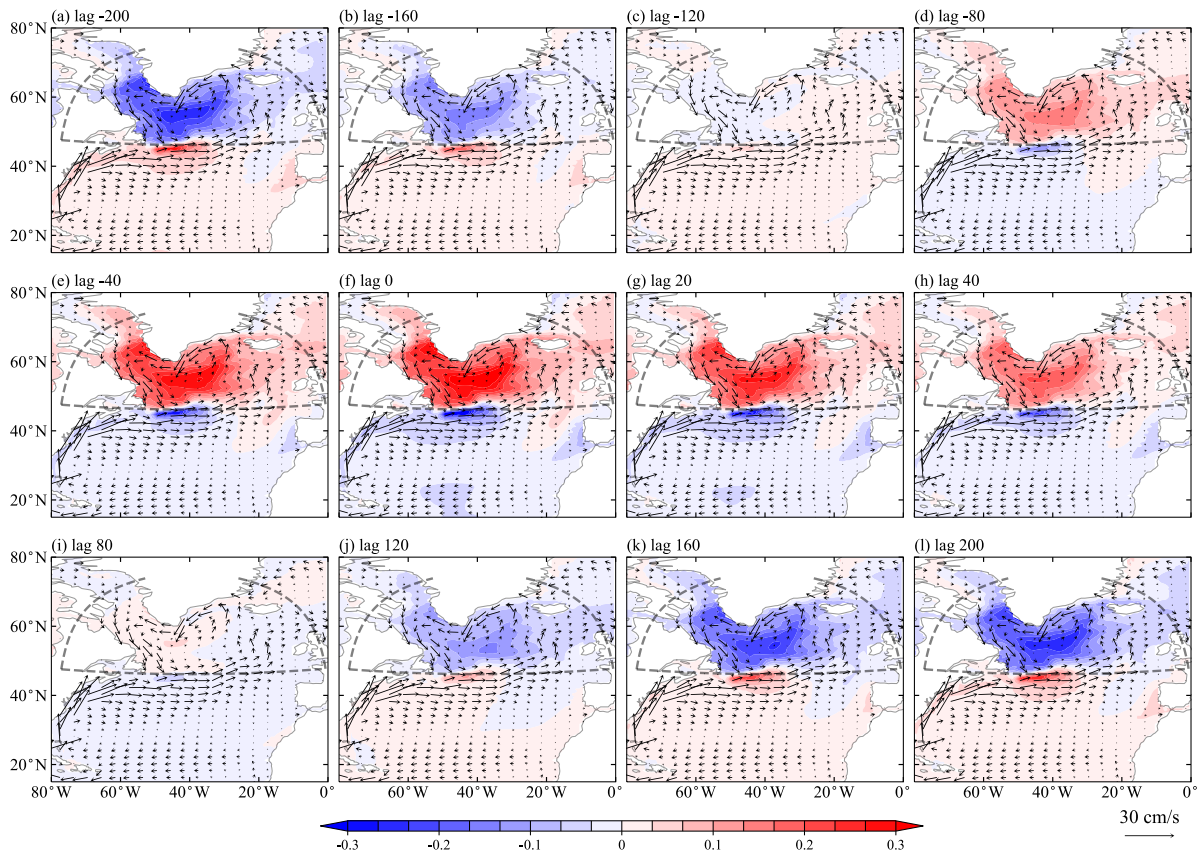
378

### 379 *b. Horizontal patterns*

380 To depict and also to explain the evolutions of salinity anomalies in these three levels, their  
 381 horizontal lead/lag regression maps on the AMOC LFC1 are plotted in Figs. 6, 8, and 10,  
 382 superimposed with climatological currents. In the North Atlantic upper ocean, the climatological  
 383 currents feature the northward Gulf Stream and its eastward extension, the northeastward NAC, and  
 384 the subpolar cyclonic circulation occupying the Labrador Sea and Irminger Sea (Fig. 6). Due to the  
 385 blocking of the “transition zone”, meridional mean currents at boundary 1 are only obvious in the  
 386 east. Mean currents at boundaries 2, 3 are significantly weaker than those at boundary 1. Therefore,  
 387 evident effect of mean advection of salinity anomaly on the DWF region is only possible at boundary  
 388 1. Even so, salinity anomaly in the DWF region evolves mostly locally without clear influence from  
 389 mean advection. For example, when salinity anomaly in the DWF region evolves from negative to  
 390 positive (Figs. 6a-f), it is always in antiphase with the subtropical salinity anomaly, and no continuous  
 391 propagation of salinity anomaly through mean advection is clearly reflected. Similarly, when the  
 392 DWF region salinity anomaly transitions from positive to negative (Figs. 6g-l), there is also little  
 393 contribution from mean advection of subtropical salinity anomaly. When salinity anomaly in the  
 394 upper DWF region is neutral (Figs. 6c, i), there appears to be weak salinity anomaly in the subtropical  
 395 upper ocean that is advected northeastward by the NAC across boundary 1, and then northwestward  
 396 by the Irminger current. However, this mean advection of salinity anomaly is too weak and may not



397 be enough to determine the weakening (Figs. 6a-c, g-i) and phase transition (Figs. 6c, i) of salinity  
 398 anomaly in the DWF region.



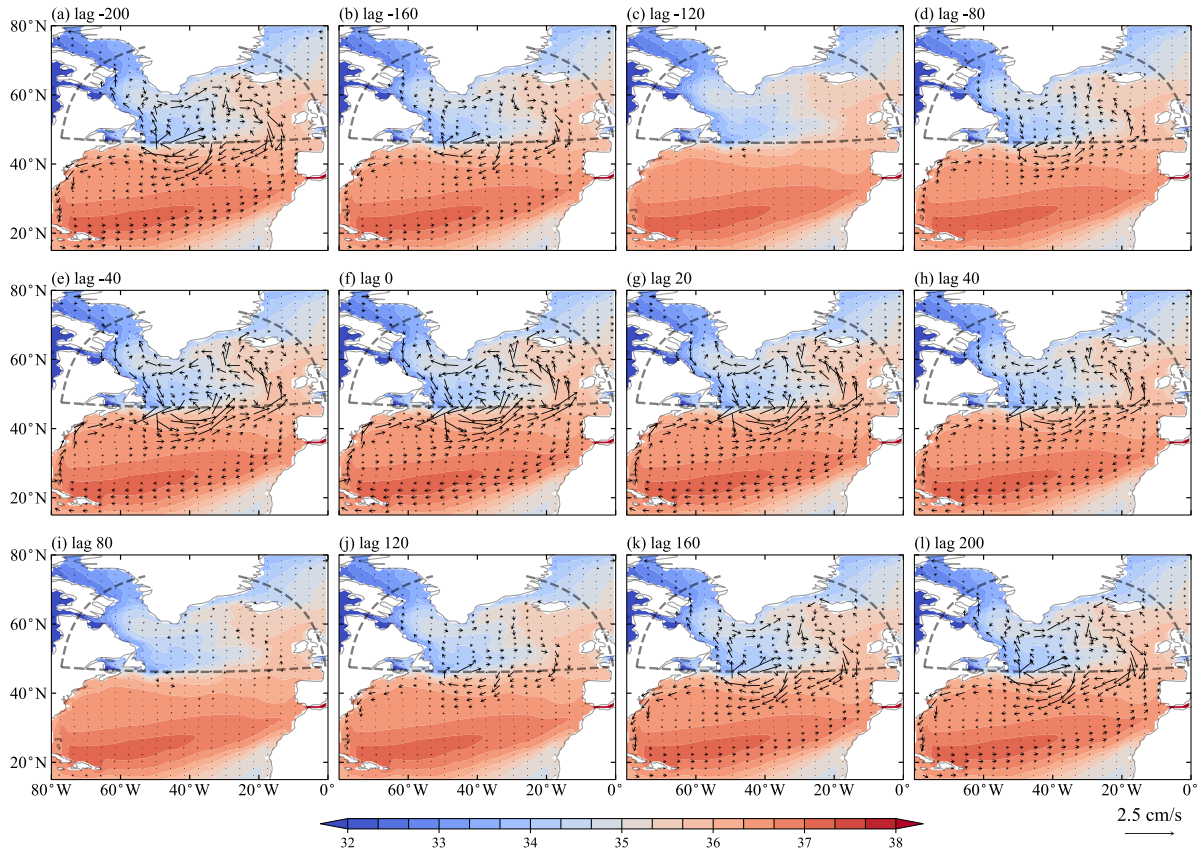
399

400 FIG. 6. Lead/lag regression coefficients of salinity anomalies averaged over 0-200 m on the AMOC LFC1  
 401 (units: psu), superimposed with climatological currents averaged over the same depth range (vector; units:  $\text{cm s}^{-1}$ ).  
 402 Negative lag means the AMOC LFC1 lags the salinity anomalies (units: year). Boundaries 1-3 defined in Fig. 2 are  
 403 also plotted.

404

405 The growth of salinity anomaly in the upper DWF region (Figs. 6d-f, j-l) is controlled by the  
 406 anomalous advection of mean salinity in the upper ocean. Figure 7 shows the lead/lag regression  
 407 coefficients of current anomalies vertically averaged over 0-200 m on the AMOC LFC1, overlaid  
 408 with climatological salinity. Current anomalies at boundaries 2, 3 are negligible throughout the entire  
 409 cycle, indicating that the higher-latitude upper ocean has little effect on the DWF region salinity  
 410 anomaly through anomalous advection of mean salinity. During the enhancing period of the positive  
 411 DWF region salinity anomaly (Figs. 6d-f), the eastward and northward NAC is also intensifying  
 412 (Figs. 7d-f), transporting more saline water from the mid-latitude eastern Atlantic to the DWF region  
 413 and enhancing the positive salinity anomaly therein. Similarly, during the enhancing period of the  
 414 negative DWF region salinity anomaly (Figs. 6j-l), the eastward and northward NAC is weakening

415 (Figs. 7j-l), reducing the northward transport of saline water from the mid-latitude eastern Atlantic  
 416 and hence enhancing the negative DWF region salinity anomaly. Therefore, salinity anomaly in the  
 417 upper DWF region and thus the AMOC anomaly are always enhanced by the anomalous advection in  
 418 the subtropical-subpolar upper ocean. This is the well-known positive salinity advection feedback  
 419 between AMOC anomaly and anomalous advection of mean salinity (Stommel 1961; Nakamura et al.  
 420 1994; Marotzke and Stone 1995; Sévellec et al. 2006).



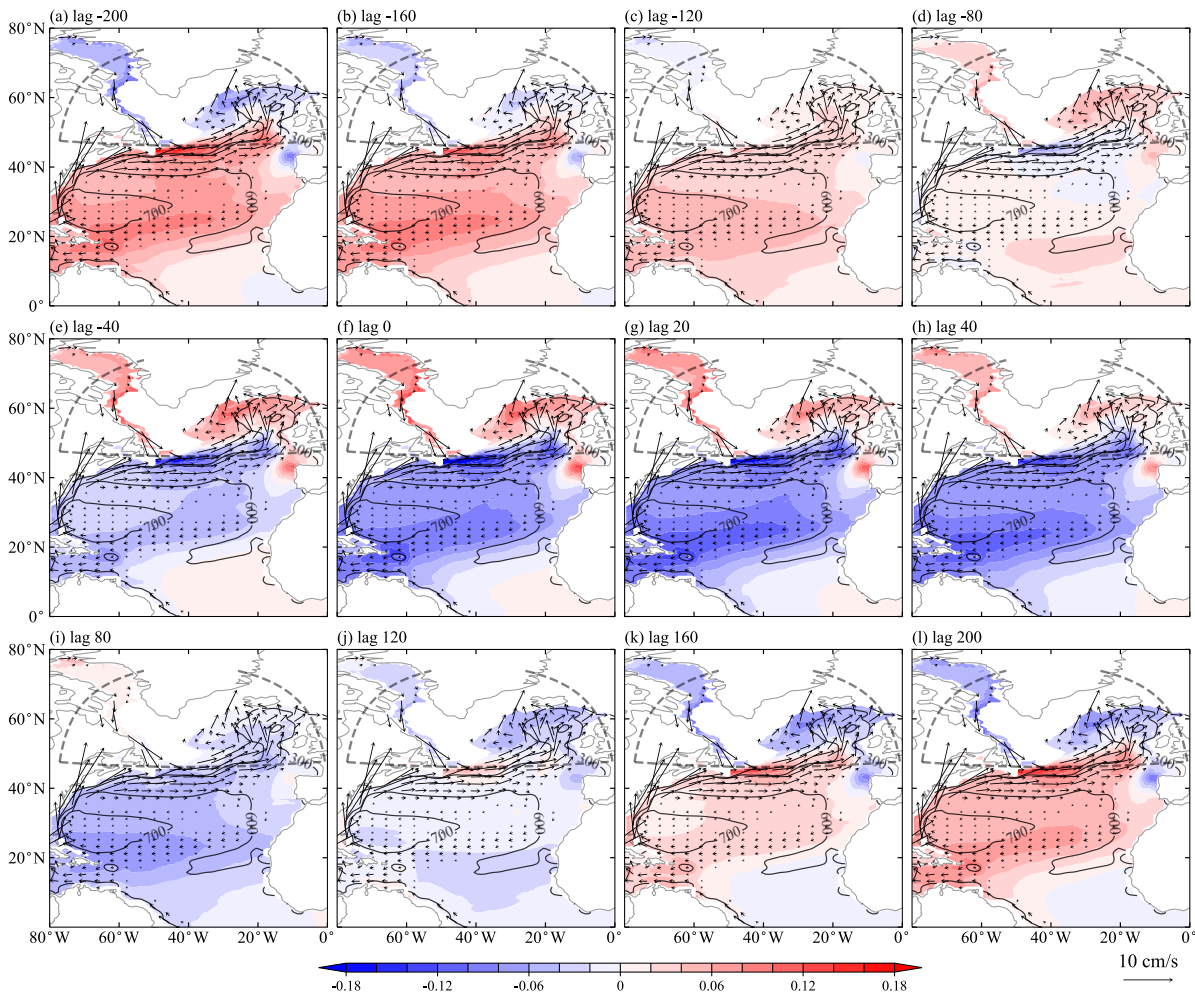
421

422 FIG. 7. Lead/lag regression coefficients of current anomalies averaged over 0-200 m on the AMOC LFC1  
 423 (units:  $\text{cm s}^{-1}$ ), superimposed with climatological salinity averaged over the same depth range (shading; units: psu).  
 424 Negative lag means the AMOC LFC1 lags the current anomalies (units: year).

425

426 It is the mean advection of salinity anomaly in the intermediate ocean that weakens the salinity  
 427 anomaly in the upper DWF region. In the intermediate North Atlantic (Fig. 8), clear mean advection  
 428 of salinity anomaly from the mid-latitude eastern Atlantic into the DWF region is reflected, having  
 429 larger influence on the DWF region than that in the upper ocean. Salinity anomalies shallower than  
 430 200 m are removed, resulting in white areas in most of the subpolar basin (Fig. 8), where the  
 431 climatological density exceeds  $27.6 \sigma_\theta$ . When salinity anomaly in the upper DWF region changes

432 from negative to positive (Figs. 6a-f), the mid-latitude salinity anomaly with the opposite sign is  
 433 advected eastward, northward, and upward by the mean NAC along isopycnals (Figs. 8a-f),  
 434 weakening salinity anomaly in the upper DWF region and contributing to its phase change. Similar  
 435 scenario also occurs when salinity anomaly in the upper DWF region changes from positive to  
 436 negative (Figs. 6g-l, 8g-l). This is the negative feedback between AMOC anomaly and mean  
 437 advection of salinity anomaly. When weakening processes outweigh enhancing processes for salinity  
 438 anomaly in the upper DWF region, its magnitude peaks and starts to decrease (Figs. 6a, f, l),  
 439 facilitating its cyclic evolution.



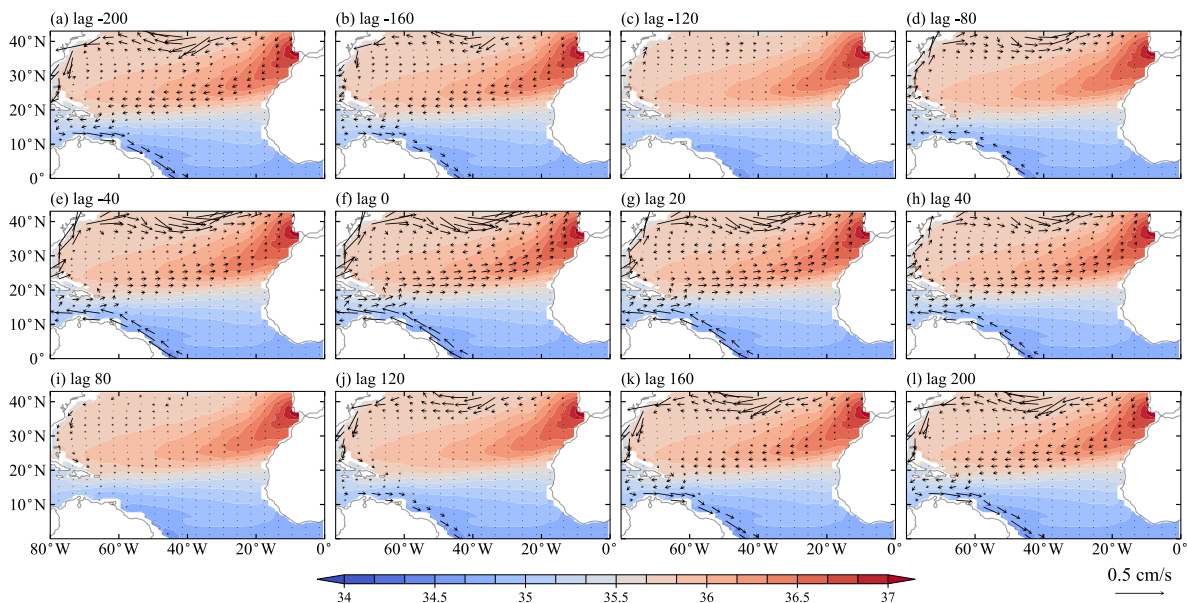
440

441 FIG. 8. Same as Fig. 6, but for variables averaged over  $26.5-27.6 \sigma_\theta$ . Black contours represent mean depths of  
 442  $26.5-27.6 \sigma_\theta$  (units: m). Currents weaker than  $0.4 \text{ cm s}^{-1}$  are not plotted. The 0-200 m salinity and current anomalies  
 443 are removed to exclude the influence from the upper ocean, resulting in the blank regions in the subpolar basin.

444

445 Both Figs. 8 and 5b depict a distinct salinity anomaly in the  $10^\circ-35^\circ\text{N}$  subtropical intermediate  
 446 ocean, which is consistently out of phase with that in the upper DWF region. Analysis reveals that,

447 this subtropical intermediate salinity anomaly is enhanced primarily by the anomalous equatorial  
 448 western boundary current (WBC), and weakened by the northward mean advection along the Gulf  
 449 Stream; surface processes as well as the vertical transport between the deep and intermediate oceans  
 450 exert only marginal effects (figure not shown). From lag -200 to -120 years (Figs. 8a-c), the mean  
 451 Gulf Stream moves positive salinity anomaly away from the subtropics into higher latitudes,  
 452 contributing to the weakening of the subtropical positive salinity anomaly. On the other hand, the  
 453 southeastward equatorial WBC anomaly reduces freshwater transport from the equatorial region into  
 454 the subtropics (Figs. 9a-c), hence enhancing the subtropical positive salinity anomaly. When the  
 455 subtropical salinity anomaly grows from slightly negative to the maximum negative value (Figs. 8e-  
 456 g), the mean Gulf Stream removes negative salinity anomaly northward, thereby again counteracting  
 457 the growth of the subtropical salinity anomaly. As the equatorial WBC anomaly turns northwestward  
 458 (Figs. 9e-g), it transports more freshwater into the subtropics and enhances the negative salinity  
 459 anomaly therein.



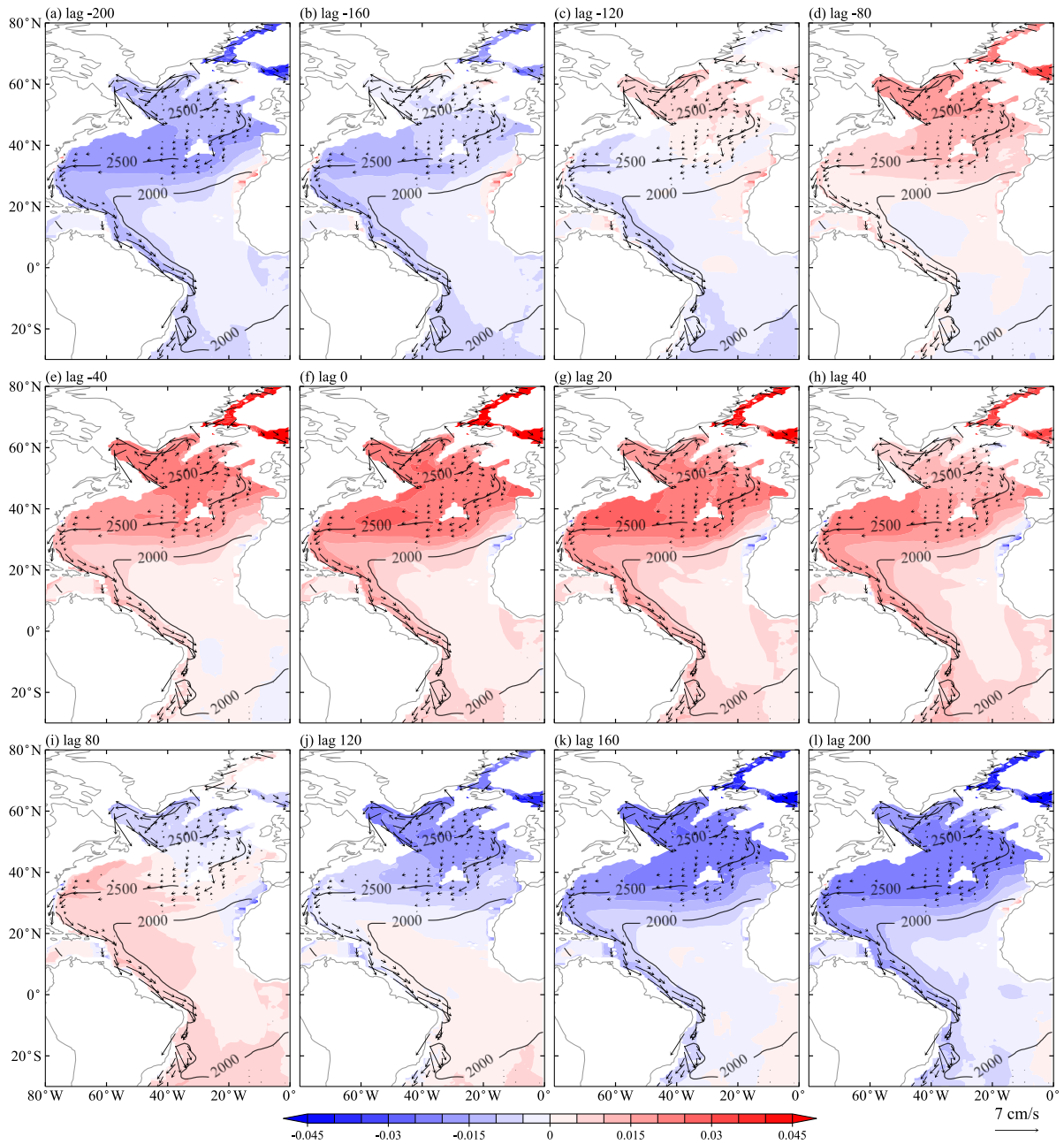
460

461 FIG. 9. Same as Fig. 7, but for variables averaged over  $26.5-27.6 \sigma_\theta$ . The 0-200 m salinity and current  
 462 anomalies are removed to exclude the influence from the upper ocean.

463

464 In the deep ocean (Fig. 10), the evolution of salinity anomalies is dominated by the mean  
 465 advection. Salinity anomalies are coherently advected southward from the subpolar basin to the  
 466 tropics and South Atlantic, through the mean southward currents especially the Deep Western  
 467 Boundary Current (DWBC). Salinity anomalies in the deep ocean exhibit nearly uniform polarity in  
 468 the whole basin, consistent with the polarity of that in the upper DWF region. The newly developed

469 positive and negative salinity anomalies in the subpolar North Atlantic (Figs. 10c, i) reach the South  
 470 Atlantic in approximately 50 years (Figs. 10d, j). Their magnitude decreases along the route, as also  
 471 observed in Fig. 5c.



472

473 FIG. 10. Same as Fig. 8, but for variables averaged over 27.75-27.85  $\sigma_\theta$ . Black contours represent mean depths  
 474 of 27.75-27.85  $\sigma_\theta$  (units: m). Currents weaker than 0.07 cm s<sup>-1</sup> are not plotted. The 0-200 m salinity and current  
 475 anomalies are removed to exclude the influence from the upper ocean, resulting in the blank regions in the subpolar  
 476 basin.

477

478 *c. Summary on collaborative salinity and AMOC evolutions*

479 Now, a three-dimensional picture of the salinity evolution in conjunction with the AMOC appears.  
480 Starting from the freshest upper DWF region and thus the weakest AMOC (Fig. 6a), the mean NAC  
481 transports positive salinity anomaly northeastward and upward from the mid-latitude intermediate  
482 ocean to the upper DWF region (Figs. 8a-c), weakening the negative salinity anomaly therein (Figs.  
483 6a-c). Meanwhile, the southward NAC anomaly on the eastern flank of boundary 1 reduces the  
484 transport of saline subtropical upper water into the DWF region, counteracting the intermediate ocean  
485 mean advection. Salinity anomaly in the subtropical intermediate ocean is positive and in antiphase  
486 with that in the upper DWF region. It is weakened by the mean Gulf Stream through losing positive  
487 salinity anomaly northward, enhanced mainly by the southeastward equatorial WBC anomaly which  
488 leads to less northward equatorial freshwater transport, and hardly affected by the deep ocean.  
489 Concurrently, negative salinity anomaly in the upper DWF region quickly descends to the deep ocean  
490 (Figs. 4a-c) and moves southward through the mean DWBC (Figs. 10a-c). These processes take about  
491 80-90 years in total, followed by phase changes of salinity anomalies in the upper DWF region and  
492 subtropical intermediate ocean (Figs. 6d, 8e). Afterward, the newly developed positive AMOC  
493 anomaly and salinity anomaly in the upper DWF region strengthen themselves through the positive  
494 salinity advection feedback, realized as the anomalous NAC's northward advection of mean salinity  
495 (Figs. 7d-f). The strengthened positive AMOC anomaly also enhances the negative salinity anomaly  
496 in the intermediate subtropics (Figs. 8e-g), through increasing the northward freshwater transport  
497 from the equatorial region (Figs. 9e-g). The positive salinity anomaly in the upper DWF region is  
498 transported downward (Figs. 4d-f) and then carried southward (Figs. 10d-f) through the mean DWBC.  
499 These processes also take about 80-90 years. Now, a half cycle of the evolutions of salinity anomalies  
500 and the AMOC is completed, taking about 180 years in total. Subsequently, the weakening processes  
501 for salinity anomaly in the upper DWF region surpass the enhancing processes, and the evolutions of  
502 salinity and AMOC anomalies enter the opposite phase.

503

## 504 **5. Determinant processes of the AMOC MCO**

### 505 *a. North Atlantic advection controls salinity anomaly in the upper DWF region*

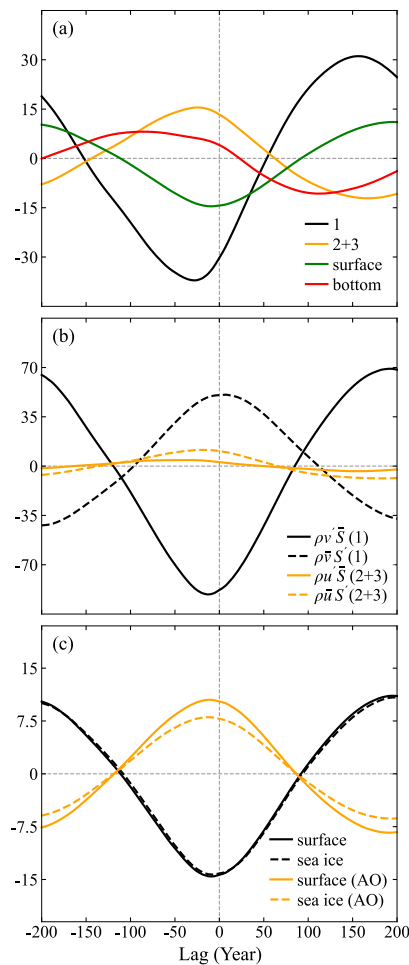
506 Salinity anomaly in the upper DWF region primarily originates from the anomalous advection of  
507 mean salinity and mean advection of salinity anomaly in the subtropical-subpolar North Atlantic. For  
508 the 0-1000 m DWF region, the freshwater mass budget comprises four components: (1) total

509 (Eulerian-mean + eddy-induced velocities) liquid freshwater mass transport at boundary 1; (2) total  
510 liquid freshwater mass transport at boundaries 2, 3; (3) surface freshwater mass flux induced by  
511 evaporation, precipitation, river runoff, and sea ice-related processes; and (4) total liquid freshwater  
512 mass transport at the 1000 m bottom. Lead/lag regression analysis of these freshwater budget  
513 components on the AMOC LFC1 is conducted in Fig. 11.  $S_{ref}$  in Eqs. (1) and (2) is given as 35.1 psu,  
514 the spatially averaged climatological salinity of the 0-1000 m DWF region.

515 Of the four aforementioned processes affecting salinity anomaly of the 0-1000 m DWF region, the  
516 most significant is the liquid freshwater transport at boundary 1 (Fig. 11a, black curve); vertical  
517 transport at the bottom is the least impactful (Fig. 11a, red curve). Freshwater transport at boundary 1  
518 diminishes freshwater input of the DWF region when the AMOC is stronger than usual, enhancing the  
519 anomalies of the DWF region salinity and AMOC. In contrast, freshwater transport at boundaries 2, 3  
520 (Fig. 11a, orange curve) weakens the AMOC anomaly, and is about half the magnitude of that of  
521 boundary 1. However, decomposition of the total liquid freshwater transport into components induced  
522 by the anomalous and mean advections in Fig. 11b reveals that, the anomalous and mean advections  
523 at boundary 1 (Fig. 11b, black curves) are about 20 and 4 times stronger than at boundaries 2, 3 (Fig.  
524 11b, orange curves), respectively. Specifically, anomalous advection of mean salinity at boundary 1  
525 (Fig. 11b, solid black curve) enhances salinity anomaly in the DWF region, manifesting the positive  
526 salinity advection feedback. Mean advection of salinity anomaly weakens salinity anomaly in the  
527 DWF region (Fig. 11b, dashed black curve), offsetting the anomalous advection of mean salinity to a  
528 large extent. As the regression coefficient of the anomalous advection-induced freshwater transport at  
529 boundaries 2, 3 shares the same sign as that of the mean advection-induced transport (Fig. 11b,  
530 orange curves), but those of boundary 1 are in antiphase with each other, Fig. 11a gives the  
531 impression that the effect of boundary 1 (Fig. 11a, black curve) on the DWF region salinity anomaly  
532 is only two times that of boundaries 2, 3 (Fig. 11a, orange curve). In fact, the actual determinants of  
533 the DWF region salinity anomaly are the anomalous and mean advections at 0-1000 m boundary 1.

534 The surface freshwater flux into the DWF region (Fig. 11c, solid black curve) is predominantly  
535 attributed to sea ice-related processes (Fig. 11c, dashed black curve). Sea ice-induced freshwater flux  
536 into the DWF region decreases when the AMOC is stronger than usual, further enhancing the AMOC  
537 anomaly. Surface freshwater flux into the Arctic Ocean region enclosed by boundaries 2, 3, and the  
538 Bering Strait, but excluding the Hudson Bay and Baltic Sea, is also evaluated (Fig. 11c, solid orange  
539 curve). During a stronger AMOC, more sea ice is melted in the Arctic Ocean (Fig. 11c, dashed orange

540 curve). On the one hand, this increases the mean advection of salinity anomaly into the DWF region  
 541 at boundaries 2, 3 (Fig. 11b, dashed orange curve). On the other hand, this reduces the sea ice  
 542 transport into the DWF region and the further sea ice melting-induced freshwater flux therein (Fig.  
 543 11c, dashed black curve). Consequently, the effects of boundaries 2, 3 on the positive (negative)  
 544 salinity anomaly in the DWF region during a stronger (weaker) AMOC are dual: weaken it through  
 545 the increased (reduced) mean advection-induced liquid freshwater import, and enhance it through the  
 546 reduced (increased) sea ice import. These two processes counteract each other, rendering the  
 547 cumulative contribution of boundaries 2, 3 to the DWF region salinity anomaly less pronounced than  
 548 when considered in isolation. This further underscores that the most decisive processes for salinity  
 549 anomaly in the 0-1000 m DWF region, and thus this North Atlantic Ocean-originated AMOC MCO,  
 550 are the anomalous advection of mean salinity and mean advection of salinity anomaly in the  
 551 subtropical-subpolar North Atlantic.



552

553 FIG. 11. Lead/lag regression coefficients of freshwater mass budget components induced by different processes  
 554 onto the AMOC LFC1 (units:  $10^6 \text{ kg s}^{-1}$ ). (a) Total liquid freshwater mass transports into the 0-1000 m DWF region



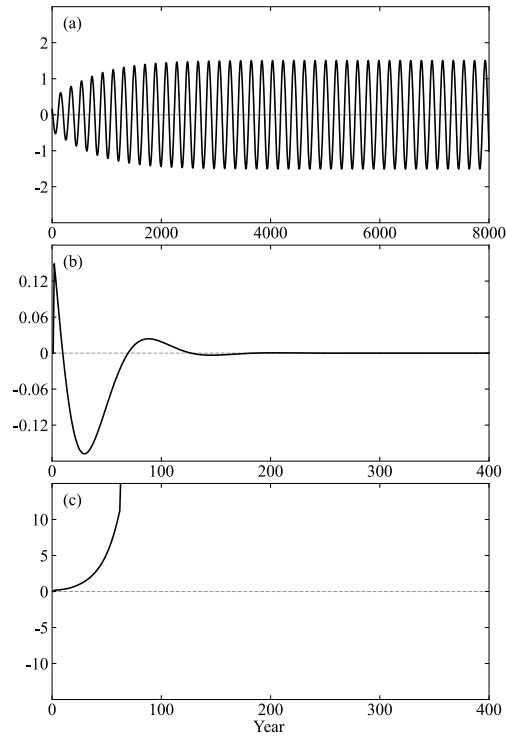
555 at boundary 1 (black curve), boundaries 2 and 3 (orange curve). Surface freshwater mass flux into the DWF region  
 556 (green curve). Total liquid freshwater mass transport at the 1000 m bottom of the DWF region (red curve). (b)  
 557 Liquid freshwater mass transports into the 0-1000 m DWF region at boundary 1, induced by the anomalous  
 558 advection of mean salinity (solid black curve) and mean advection of salinity anomaly (dashed black curve). The  
 559 orange curves are the same as the black curves, but for boundaries 2, 3. (c) Surface freshwater mass flux into the  
 560 DWF region (solid black curve) and its sea ice-induced component (dashed black curve). The orange curves are the  
 561 same as the black curves, but for the Arctic Ocean region encircled by boundaries 2, 3, and the Bering Strait (not  
 562 annotated in figures), but excluding the Hudson Bay and Baltic Sea. Negative lag means the AMOC LFC1 lags the  
 563 freshwater terms (units: year).

564

### 565 *b. Indispensability of the North Atlantic anomalous and mean advectons*

566 Utilizing the North Atlantic-only box model in LY22 (appendix), we will examine whether this  
 567 North Atlantic Ocean-originated AMOC MCO can occur when either the anomalous advection or the  
 568 mean advection in the AMOC's upper limb is artificially deactivated. The theoretical model does not  
 569 distinguish between the upper and intermediate oceans. In the upper branch, the anomalous advection  
 570 of mean salinity and mean advection of salinity anomaly are represented by  $q'(\overline{S}_1 - \overline{S}_2)$  and  
 571  $\overline{q}(S'_1 - S'_2)$ , respectively. Starting with a positive  $q'$ ,  $q'(\overline{S}_1 - \overline{S}_2)$  increases  $S'_2$  and further strengthens  
 572  $q'$ , constituting the positive salinity advection feedback. Concurrently,  $q'(\overline{S}_4 - \overline{S}_1)$  renders  $S'_1$   
 573 negative through increasing freshwater transport from subtropical deeper box 4 into subtropical upper  
 574 box 1. The negative  $S'_1$  is then advected northward through  $\overline{q}(S'_1 - S'_2)$ , weakening the positive  $S'_2$  and  
 575 thus the positive  $q'$ . This constitutes the mean advection process.  $q'(\overline{S}_4 - \overline{S}_1)$  represents a difference  
 576 between the CESM1 simulation and the theoretical model. In the CESM1 simulation, the major  
 577 anomalous advection that affects salinity anomaly in the subtropics is horizontal and originated from  
 578 the equatorial region (Fig. 9). However, the paramount processes contributing to the DWF region  
 579 salinity anomaly: anomalous and mean advectons in the subtropical-subpolar North Atlantic, are  
 580 present in both the CESM1 simulation and theoretical model.

581 When Eq. (A) is active, a self-sustained AMOC oscillation is exhibited in the theoretical model  
 582 (Fig. 12a). When artificially deactivating the anomalous advection in the upper branch [ $q'(\overline{S}_1 - \overline{S}_2)$ ]  
 583 while leaving other processes unmodified, the AMOC exhibits a strongly damped oscillation (Fig.  
 584 12b), because of the absence of the positive salinity advection feedback that enhances the AMOC  
 585 anomaly. On the other hand, when the mean advection in the upper branch [ $\overline{q}(S'_1 - S'_2)$ ] is  
 586 deactivated, its weakening effect on the AMOC anomaly is eliminated, therefore leading to a runaway  
 587 tendency for  $q'$  (Fig. 12c). This further reveals the indispensability of the anomalous and mean  
 588 advectons in the upper AMOC limb for the North Atlantic Ocean-originated AMOC MCO.



589

590 FIG. 12. (a) Self-sustained oscillation of the AMOC anomaly  $q'$  (units: Sv) in the theoretical model (appendix).  
 591 (b) Damped oscillation of  $q'$  when the anomalous advection of mean salinity  $q'(\overline{S}_1 - \overline{S}_2)$  in the upper branch is  
 592 deactivated. (c) Runaway  $q'$  when the mean advection of salinity anomaly  $\overline{q}(S'_1 - S'_2)$  in the upper branch is  
 593 deactivated.

594

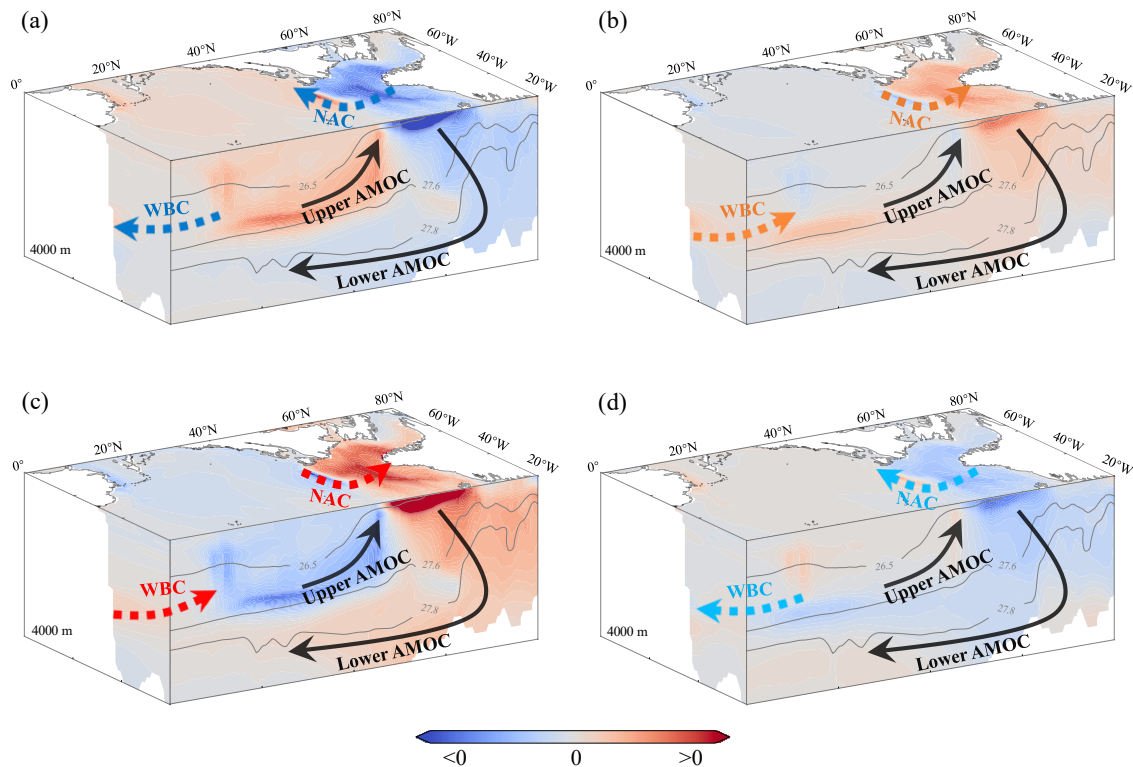
## 595 6. Summary and discussion

596 An MCV of the AMOC is identified in a CESM1 control simulation. It is interpreted as a linear  
 597 oscillation and can be termed an MCO. This AMOC MCO is primarily driven by processes in the  
 598 North Atlantic, thereby potentially representing a North Atlantic Ocean-originated mode of the real-  
 599 world AMOC MCV. In the upper level of the DWF region in the subpolar North Atlantic, variation of  
 600 salinity anomaly dominates variation of density anomaly, leading to the AMOC oscillation. The most  
 601 determinant processes for salinity anomaly in the upper DWF region, and thus the AMOC MCO, are  
 602 the anomalous advection of mean salinity and mean advection of salinity anomaly in the subtropical-  
 603 subpolar North Atlantic (Fig. 11b, black curves), which have much greater effects than other  
 604 processes. Contribution from the bottom of the DWF region is negligible. The Arctic Ocean exerts  
 605 two counteractive effects on this AMOC MCO. First, a positive AMOC anomaly increases Arctic  
 606 Ocean sea ice melting, resulting in a negative salinity anomaly therein that can be advected into the  
 607 DWF region through mean advection. This curbs the AMOC anomaly. Second, as more sea ice melts

608 in the Arctic Ocean, there is a reduction in sea ice transport into the DWF region, thus decreasing the  
609 melting-induced freshwater input of the DWF region. This enhances the AMOC anomaly. Overall,  
610 the cumulative effect of the Arctic Ocean on this AMOC MCO is relatively minor, and much weaker  
611 than the anomalous and mean advections in the North Atlantic. Additionally, no obvious salinity  
612 anomaly that evolves synchronously with the AMOC is observed in the South Atlantic or Southern  
613 Ocean. Therefore, it is appropriate to designate this AMOC MCO as a North Atlantic Ocean-  
614 originated AMOC MCO.

615 Figure 13 schematically summarizes the core processes for this AMOC MCO, with an emphasis  
616 on advection in the North Atlantic. At the AMOC's negative peak (Fig. 13a), both the NAC and  
617 equatorial WBC display southward anomalies (Fig. 13a, dashed blue arrows). The negative salinity  
618 anomaly in the upper DWF region is enhanced by the southward NAC anomaly (Fig. 13a, upper  
619 dashed blue arrow) through the reduced northward subtropical saline water transport, and weakened  
620 by the northward mean advection (Fig. 13, upper solid black arrow) that carries the positive salinity  
621 anomaly in the subtropical intermediate ocean northward. Salinity anomaly in the upper DWF region  
622 constantly descends through the subpolar convection or vertical mixing, then moves southward  
623 through the mean DWBC (Fig. 13, lower solid black arrow). The positive salinity anomaly in the  
624 subtropical intermediate ocean is weakened by the northward mean advection that carries salinity  
625 anomaly away northward, and enhanced by the reduced northward freshwater transport induced by  
626 the southward equatorial WBC anomaly (Fig. 13a, lower dashed blue arrow). At this stage,  
627 weakening processes for the negative (positive) salinity anomaly in the upper DWF region  
628 (subtropical intermediate ocean) are stronger than enhancing processes, hence the negative DWF  
629 region (positive subtropical intermediate) salinity anomaly is weakening.

630 After the phase transition (Fig. 13b), anomalies of the upper DWF region salinity and AMOC turn  
631 slightly positive, as do the NAC and equatorial WBC (Fig. 13b, dashed orange arrows). These  
632 anomalies are subsequently strengthened by the positive salinity advection feedback. The subtropical  
633 intermediate salinity anomaly lags slightly behind that in the upper DWF region, and is still  
634 undergoing its phase transition. Later, it turns negative and is strengthened by the increased northward  
635 freshwater transport through the northward equatorial WBC anomaly. When enhancing processes for  
636 salinity anomaly in the upper DWF region are surpassed by weakening processes, the positive AMOC  
637 anomaly peaks (Fig. 13c) and starts to neutralize. Subsequently, the AMOC anomaly turns negative  
638 (Fig. 13d) and progresses toward the negative peak (Fig. 13a), completing a full cycle.



639

640 FIG. 13. Schematic diagrams showing oceanic states in the North Atlantic during (a) the peak of the negative  
 641 phase, (b) the start of the positive phase, (c) the peak of the positive phase, and (d) the start of the negative phase of  
 642 the AMOC multacentennial oscillation (MCO). Dashed and solid arrows represent anomalous and mean currents,  
 643 respectively.

644

645 Significance of the two paramount processes for this AMOC MCO: the anomalous and mean  
 646 advectons in the upper AMOC branch, is tested utilizing the theoretical model in LY22. When both  
 647 processes are active, the theoretical model exhibits a pronounced AMOC MCO. Equations of the  
 648 theoretical model [Eq. (A)], which capture the essence of this modeled North Atlantic Ocean-  
 649 originated AMOC MCO, can mathematically explain the enhancing and weakening effects of the  
 650 anomalous and mean advectons on the AMOC anomaly. If either the anomalous or mean advection  
 651 in the upper branch is artificially deactivated, the theoretical model cannot exhibit the AMOC MCO,  
 652 emphasizing the essential role of these two advection processes in this North Atlantic Ocean-  
 653 originated AMOC MCO.

654 The AMOC MCO analyzed in this study is North Atlantic Ocean-originated, differentiating it  
 655 from the “flip-flop” AMOC MCO in Park and Latif (2008), the Southern Ocean-originated AMOC

656 MCO in Delworth and Zeng (2012), as well as the Arctic Ocean-originated AMOC MCO in Jiang et  
657 al. (2021) and Meccia et al. (2023). The “flip-flop” AMOC MCO represents a multi-equilibrium  
658 phenomenon that is markedly distinct from our study. The main difference between the Southern  
659 Ocean-originated AMOC MCO and our North Atlantic Ocean-originated AMOC MCO lies in the  
660 location of the salinity anomaly that is advected northward toward the DWF region by mean  
661 advection. In their study, mean advection moves salinity anomaly in the upper Southern Ocean  
662 northward, whereas in ours, the northward salinity anomaly originates from the subtropical  
663 intermediate ocean. In Jiang et al. (2021) and Meccia et al. (2023), clear current and salinity  
664 anomalies are exhibited north of the subpolar North Atlantic, yet in our study the anomalies therein  
665 are rather weak, especially the current anomalies. Given the connection between salinity anomalies  
666 from the Arctic Ocean and the sea ice thermodynamics, distinctions between the Arctic Ocean-  
667 originated and North Atlantic Ocean-originated AMOC MCOs likely stem from the difference in both  
668 the ocean model and the sea ice model utilized.

669 Despite these differences, our study still aids the understanding of other AMOC MCOs. The  
670 positive salinity advection feedback in the subtropical-subpolar upper ocean is pivotal not only in this  
671 study, but also in the Southern Ocean-originated and Arctic Ocean-originated AMOC MCOs. In Fig.  
672 5a of Delworth and Zeng (2012), the continuous northward salinity anomaly that is symbolic of mean  
673 advection reaches approximately 45°N. North of 45°N, salinity anomalies evolve nearly  
674 synchronously, mirroring the corresponding pattern in Fig. 5a of our current study. Therefore, this  
675 local evolution of salinity anomaly is likely driven by the positive salinity advection feedback. Local  
676 salinity evolution north of 45°N is also reflected in Fig. 5a in Jiang et al. (2021) and Fig. 5a in Meccia  
677 et al. (2023), suggesting that the positive salinity advection feedback is also likely to be the enhancing  
678 process for AMOC anomaly in the Arctic Ocean-originated AMOC MCOs, which is not addressed in  
679 these two studies. The AMOC MCO identified in the intermediate-complexity model study of  
680 Mehling et al. (2023) has a similar mechanism to the Arctic Ocean-originated AMOC MCOs. They  
681 employed a box model adapted from Stommel (1961) to explain their AMOC MCO, and highlighted  
682 the Arctic Ocean-originated salinity anomaly by incorporating an additional Arctic Ocean box.  
683 Likewise, Wei and Zhang (2022) also utilized a revised Stommel’s two-box model incorporating a  
684 negative feedback representing salinity anomaly from the Arctic Ocean, to account for the Arctic  
685 Ocean-originated AMOC multidecadal oscillation. Its essence aligns closely with that of the Arctic  
686 Ocean-originated AMOC MCOs. Both theoretical models employed in these two studies actually

687 incorporate the positive salinity advection feedback and mean advection process in the subtropical-  
688 subpolar upper ocean, although their focus is salinity anomaly from the Arctic Ocean. Hence, the  
689 subtropical-subpolar positive salinity advection feedback likely serves as the essential enhancing  
690 process for AMOC anomaly in the North Atlantic Ocean-originated, Southern Ocean-originated, and  
691 Arctic Ocean-originated AMOC MCOs. The primary difference among them is perhaps the origin of  
692 the salinity anomaly that is advected into the DWF region through mean advection. With the  
693 incorporation of additional boxes representing the South Atlantic/Southern Ocean, the theoretical  
694 model in LY22 has the potential to account for the South Ocean-originated AMOC MCO, through  
695 capturing salinity anomalies in the Southern Ocean.

696 To the best of our knowledge, this study is perhaps the first to identify the North Atlantic Ocean-  
697 originated AMOC MCO in coupled models. Therefore, this mode will be more convincing if  
698 identified in other coupled models. However, the assessments of whether this North Atlantic Ocean-  
699 originated AMOC MCO and other previously proposed AMOC MCOs genuinely exist in the Earth's  
700 climate system, as well as their relative contributions to the real-world AMOC MCV, are inhibited by  
701 the limited direct observations that are unfortunately unavailable in the foreseeable future. Given that  
702 the natural forcing remains constant in these coupled model simulations, these coupled model AMOC  
703 MCOs all represent internal variability. Therefore, the AMOC MCV might have been a background  
704 for the anthropogenic centennial climate change. Further research into the mechanisms of various  
705 potential modes that constitute the real-world AMOC MCV will enhance our understanding of the  
706 ongoing climate change.

707

#### 708 *Acknowledgements.*

709 This research is jointly supported by the NSF of China (Nos. 42230403, 42288101, and  
710 41725021) and by the foundation at the Shanghai Frontiers Science Centre of Atmosphere-Ocean  
711 Interaction of Fudan University. Prof. Q. Zhang acknowledges supports from Swedish Research  
712 Council (2022-03129).

713

#### 714 *Data Availability Statement.*

715 All data used in this study are available upon request.

716

## APPENDIX

717

718 In Li and Yang (2022), we proposed a North Atlantic-only theoretical model for AMOC  
 719 multcentennial oscillation. The North Atlantic is divided into 4 boxes (Fig. A1): subtropical upper  
 720 ocean box 1, subpolar upper ocean box 2, subpolar deeper ocean box 3, and subtropical deeper ocean  
 721 box 4. Equilibrium (anomalous) salinities of boxes 1-4 are represented by  $\overline{S}_1$  ( $S'_1$ ),  $\overline{S}_2$  ( $S'_2$ ),  $\overline{S}_3$  ( $S'_3$ ),  
 722 and  $\overline{S}_4$  ( $S'_4$ ), respectively. Climatological (anomalous) AMOC is represented by  $\overline{q}$  ( $q'$ ). Only salinity  
 723 variation is considered here. The linearized salinity equations are:

$$724 \quad V_1 \dot{S}'_1 = q'(\overline{S}_4 - \overline{S}_1) + \overline{q}(S'_4 - S'_1) \quad (\text{A1})$$

$$725 \quad V_2 \dot{S}'_2 = q'(\overline{S}_1 - \overline{S}_2) + \overline{q}(S'_1 - S'_2) - k_m(S'_2 - S'_3) \quad (\text{A2})$$

$$726 \quad V_3 \dot{S}'_3 = \overline{q}(S'_2 - S'_3) + k_m(S'_2 - S'_3) \quad (\text{A3})$$

$$727 \quad V_4 \dot{S}'_4 = \overline{q}(S'_3 - S'_4) \quad (\text{A4})$$

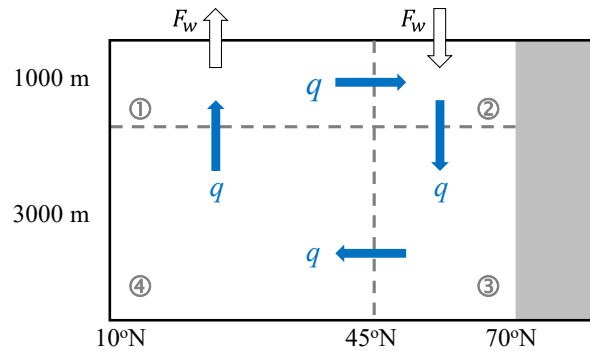
$$728 \quad q' = \lambda \Delta \rho' = \lambda \rho_0 \beta [\delta(S'_2 - S'_1) + (1 - \delta)(S'_3 - S'_4)] \quad (\text{A5})$$

$$729 \quad \delta = \frac{V_1}{V_1 + V_4} = \frac{V_2}{V_2 + V_3} = \frac{D_1}{D} \quad (\text{A6})$$

730 where  $V_1$ ,  $V_2$ ,  $V_3$ , and  $V_4$  are the volumes of boxes 1-4, respectively.  $\lambda$  is a linear closure coefficient,  
 731 representing the sensitivity of AMOC anomaly to meridional density difference  $\Delta \rho'$ .  $\rho_0$  and  $\beta$  are the  
 732 reference sea water density and haline contraction coefficient.  $k_m$  is a nonlinear vertical mixing term  
 733 parameterized as  $\kappa q'^2$ , with  $\kappa = 10^{-4} \text{ m}^{-3} \text{ s}$ .

734 In this study, the box model's geometry and parameters are set according to the CESM1 control  
 735 simulation. Boxes 1-4 span the  $10^\circ$ - $45^\circ$ N and 0-1000 m,  $45^\circ$ - $70^\circ$ N and 0-1000 m,  $45^\circ$ - $70^\circ$ N and  
 736 1000-4000 m, and  $10^\circ$ - $45^\circ$ N and 1000-4000 m of the North Atlantic domain in the CESM1 model,  
 737 respectively. Their volumes  $V_{1-4}$  are  $2.65 \times 10^{16} \text{ m}^3$ ,  $0.77 \times 10^{16} \text{ m}^3$ ,  $1.16 \times 10^{16} \text{ m}^3$ , and  $6.3 \times$   
 738  $10^{16} \text{ m}^3$ , respectively.  $\overline{q}$  is set to 24 Sv, corresponding to the climatological AMOC (Fig. 1a).  $\overline{S}_{1-4}$  are  
 739 the climatological salinities of the CESM1 model domains corresponding to boxes 1-4.  $\overline{S}_1$  is 35.9 psu.  
 740 The actual values of  $\overline{S}_{2-4}$  are all close to 35.2 psu, so they are collectively set to 35.2 psu for  
 741 conciseness.  $\rho_0$  and  $\beta$  are set to  $10^3 \text{ kg m}^{-3}$  and  $7.61 \times 10^{-4} \text{ psu}^{-1}$ , respectively. Figure A2 suggests  
 742 that in the CESM1 simulation, the AMOC anomaly is linearly proportional to the anomaly of

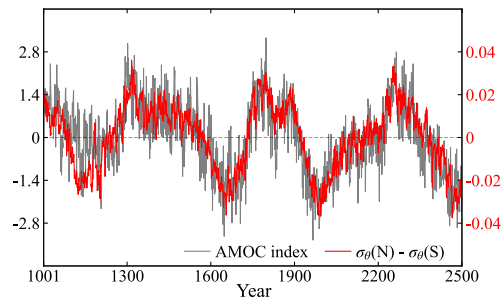
743 difference in potential density between  $45^{\circ}$ - $70^{\circ}$ N and  $10^{\circ}$ - $45^{\circ}$ N North Atlantic. According to Fig. A2,  
 744  $\lambda$  is set to  $70 \text{ Sv kg}^{-1} \text{ m}^3$ .



745

746 FIG. A1. Schematic of the 4-box ocean model. Ocean boxes are denoted by ①, ②, ③, and ④. Boxes 1 and 4  
 747 represent the upper and deeper subtropical oceans, respectively; boxes 2 and 3 represent the upper and deeper  
 748 subpolar oceans, respectively.  $D_1$  and  $D_2$  are the depths of the upper and deeper boxes, respectively.  $F_w$  is the net  
 749 freshwater flux out of (into) the subtropical (subpolar) surface ocean, which is cancelled out during the linearization  
 750 to derive Eq. (A).  $q$  represents the AMOC strength.

751



752

753 FIG. A2. Time series for the anomalies of AMOC index (gray curve, left y-axis; units: Sv) and difference in  
 754 potential density  $\sigma_{\theta}$  between the  $45^{\circ}$ - $70^{\circ}$ N and  $10^{\circ}$ - $45^{\circ}$ N North Atlantic (red curve, right y-axis; units:  $\text{kg m}^{-3}$ ), in  
 755 the CESM1 control simulation.

756



## REFERENCES

- 757  
758 Askjær, T. G., and Coauthors, 2022: Multi-centennial Holocene climate variability in proxy records  
759 and transient model simulations. *Quat. Sci. Rev.*, **296**, 20,  
760 <https://doi.org/10.1016/j.quascirev.2022.107801>
- 761 Ayache, M., D. Swingedouw, Y. Mary, F. Eynaud, and C. Colin, 2018: Multi-centennial variability of  
762 the AMOC over the Holocene: A new reconstruction based on multiple proxy-derived SST  
763 records. *Global Planet. Change*, **170**, 172-189, <https://doi.org/10.1016/j.gloplacha.2018.08.016>
- 764 Buckley, M. W., and J. Marshall, 2016: Observations, inferences, and mechanisms of the Atlantic  
765 Meridional Overturning Circulation: A review. *Rev. Geophys.*, **54**, 5-63,  
766 <https://doi.org/10.1002/2015RG000493>
- 767 Chabaud, L., M. F. S. Goni, S. Desprat, and L. Rossignol, 2014: Land-sea climatic variability in the  
768 eastern North Atlantic subtropical region over the last 14,200 years: Atmospheric and oceanic  
769 processes at different timescales. *Holocene*, **24**, 787-797,  
770 <https://doi.org/10.1177/0959683614530439>
- 771 Chapman, M. R., and N. J. Shackleton, 2000: Evidence of 550-year and 1000-year cyclicities in North  
772 Atlantic circulation patterns during the Holocene. *Holocene*, **10**, 287-291,  
773 <https://doi.org/10.1191/095968300671253196>
- 774 Craig, A. P., M. Vertenstein, and R. Jacob, 2012: A new flexible coupler for earth system modeling  
775 developed for CCSM4 and CESM1. *Int. J. High Perform. Comput. Appl.*, **26**, 31-42,  
776 <https://doi.org/10.1177/1094342011428141>
- 777 Delworth, T. L., and F. R. Zeng, 2012: Multicentennial variability of the Atlantic meridional  
778 overturning circulation and its climatic influence in a 4000 year simulation of the GFDL CM2.1  
779 climate model. *Geophys. Res. Lett.*, **39**, <https://doi.org/10.1029/2012GL052107>
- 780 Drijfhout, S., C. Heinze, M. Latif, and E. MaierReimer, 1996: Mean circulation and internal  
781 variability in an ocean primitive equation model. *J. Phys. Oceanogr.*, **26**, 559-580,  
782 [https://doi.org/10.1175/1520-0485\(1996\)026<0559:MCAIVI>2.0.CO;2](https://doi.org/10.1175/1520-0485(1996)026<0559:MCAIVI>2.0.CO;2)
- 783 Fox-Kemper, B., and R. Ferrari, 2008: Parameterization of mixed layer eddies. Part II: Prognosis and  
784 impact. *J. Phys. Oceanogr.*, **38**, 1166-1179, <https://doi.org/10.1175/2007JPO3788.1>
- 785 Fox-Kemper, B., R. Ferrari, and R. Hallberg, 2008: Parameterization of mixed layer eddies. Part I:  
786 Theory and diagnosis. *J. Phys. Oceanogr.*, **38**, 1145-1165,  
787 <https://doi.org/10.1175/2007JPO3792.1>
- 788 Gent, P. R., and J. C. McWilliams, 1990: Isopycnal mixing in ocean circulation models. *J. Phys.*

- 789 *Oceanogr.*, **20**, 150-155, [https://doi.org/10.1175/1520-0485\(1990\)020<0150:IMIOCM>2.0.CO;2](https://doi.org/10.1175/1520-0485(1990)020<0150:IMIOCM>2.0.CO;2)
- 790 Griffies, S. M., and E. Tziperman, 1995: A linear thermohaline oscillator driven by stochastic  
791 atmospheric forcing. *J. Climate*, **8**, 2440-2453, [https://doi.org/10.1175/1520-0442\(1995\)008<2440:ALTODB>2.0.CO;2](https://doi.org/10.1175/1520-0442(1995)008<2440:ALTODB>2.0.CO;2)
- 792
- 793 Hall, I. R., G. G. Bianchi, and J. R. Evans, 2004: Centennial to millennial scale Holocene climate-  
794 deep water linkage in the North Atlantic. *Quat. Sci. Rev.*, **23**, 1529-1536,  
795 <https://doi.org/10.1016/j.quascirev.2004.04.004>
- 796 Hansen, B., S. Østerhus, H. Hátún, R. Kristiansen, and K. M. H. Larsen, 2003: The Iceland-Faroe  
797 inflow of Atlantic water to the Nordic Seas. *Prog. Oceanogr.*, **59**, 443-474,  
798 <https://doi.org/10.1016/j.pocean.2003.10.003>
- 799 Hughes, Tertia M. C., and Andrew J. Weaver, 1994: Multiple equilibria of an asymmetric two-basin  
800 ocean model. *J. Phys. Oceanogr.*, **24**, 619-637, [https://doi.org/10.1175/1520-0485\(1994\)024<0619:MEOAAT>2.0.CO;2](https://doi.org/10.1175/1520-0485(1994)024<0619:MEOAAT>2.0.CO;2)
- 801
- 802 Hunke, E. C., and W. H. Lipscomb, 2010: CICE: The Los Alamos Sea Ice Model, documentation and  
803 software user's manual, version 4.1. Los Alamos National Laboratory Rep. LACC-06-012, 76  
804 pp.
- 805 Jiang, W. M., G. Gastineau, and F. Codron, 2021: Multicentennial variability driven by salinity  
806 exchanges between the Atlantic and the Arctic Ocean in a coupled climate model. *J. Adv. Model. Earth Syst.*, **13**, e2020MS002366, <https://doi.org/10.1029/2020MS002366>
- 807
- 808 Jin, F. F., 1997: An equatorial ocean recharge paradigm for ENSO. Part I: Conceptual model. *J. Atmos. Sci.*, **54**, 811-829, [https://doi.org/10.1175/1520-0469\(1997\)054<0811:AEORPF>2.0.CO;2](https://doi.org/10.1175/1520-0469(1997)054<0811:AEORPF>2.0.CO;2)
- 809
- 810
- 811 Kieke, D., and I. Yashayaev, 2015: Studies of Labrador Sea Water formation and variability in the  
812 subpolar North Atlantic in the light of international partnership and collaboration. *Prog. Oceanogr.*, **132**, 220-232, <https://doi.org/10.1016/j.pocean.2014.12.010>
- 813
- 814 Kim, J. H., and Coauthors, 2004: North Pacific and North Atlantic sea-surface temperature variability  
815 during the Holocene. *Quat. Sci. Rev.*, **23**, 2141-2154,  
816 <https://doi.org/10.1016/j.quascirev.2004.08.010>
- 817
- 818 Kissel, C., A. Van Toer, C. Laj, E. Cortijo, and E. Michel, 2013: Variations in the strength of the North  
819 Atlantic bottom water during Holocene. *Earth Planet. Sci. Lett.*, **369**, 248-259,  
820 <https://doi.org/10.1016/j.epsl.2013.03.042>
- 820 Lawrence, D. M., and Coauthors, 2011: Parameterization improvements and functional and structural

- 821 advances in version 4 of the Community Land Model. *J. Adv. Model. Earth Syst.*, **3**, 27,  
822 <https://doi.org/10.1029/2011MS000045>
- 823 Li, Y., and H. J. Yang, 2022: A theory for self-sustained multicentennial oscillation of the Atlantic  
824 meridional overturning circulation. *J. Climate*, **35**, 5883-5896, [https://doi.org/10.1175/JCLI-D-  
825 21-0685.1](https://doi.org/10.1175/JCLI-D-21-0685.1)
- 826 Liu, Z. Y., 2012: Dynamics of interdecadal climate variability: A historical perspective. *J. Climate*, **25**,  
827 1963-1995, <https://doi.org/10.1175/2011JCLI3980.1>
- 828 Marotzke, J., 1996: Analysis of thermohaline feedbacks. *Decadal Climate Variability: Dynamics and  
829 Predictability*, D. L. T. Anderson, and J. Willebrand, Eds., Springer, 333-378.
- 830 Marotzke, J., and P. H. Stone, 1995: Atmospheric transports, the thermohaline circulation, and flux  
831 adjustments in a simple coupled model. *J. Phys. Oceanogr.*, **25**, 1350-1364,  
832 [https://doi.org/10.1175/1520-0485\(1995\)025<1350:ATTTCA>2.0.CO;2](https://doi.org/10.1175/1520-0485(1995)025<1350:ATTTCA>2.0.CO;2)
- 833 Martin, T., W. Park, and M. Latif, 2013: Multi-centennial variability controlled by Southern Ocean  
834 convection in the Kiel Climate Model. *Climate Dyn.*, **40**, 2005-2022,  
835 <https://doi.org/10.1007/s00382-012-1586-7>
- 836 —, 2015: Southern Ocean forcing of the North Atlantic at multi-centennial time scales in the Kiel  
837 Climate Model. *Deep-Sea Res. Part II-Top. Stud. Oceanogr.*, **114**, 39-48,  
838 <https://doi.org/10.1016/j.dsr2.2014.01.018>
- 839 McDermott, F., D. P. Matthey, and C. Hawkesworth, 2001: Centennial-scale Holocene climate  
840 variability revealed by a high-resolution speleothem delta 18O record from SW Ireland. *Science*,  
841 **294**, 1328-1331, <https://doi.org/10.1126/science.1063678>
- 842 Meccia, V. L., R. Fuentes-Franco, P. Davini, K. Bellomo, F. Fabiano, S. T. Yang, and J. von  
843 Hardenberg, 2023: Internal multi-centennial variability of the Atlantic meridional overturning  
844 circulation simulated by EC-Earth3. *Climate Dyn.*, **18**, [https://doi.org/10.1007/s00382-022-  
845 06534-4](https://doi.org/10.1007/s00382-022-06534-4)
- 846 Mehling, O., K. Bellomo, M. Angeloni, C. Pasquero, and J. von Hardenberg, 2023: High-latitude  
847 precipitation as a driver of multicentennial variability of the AMOC in a climate model of  
848 intermediate complexity. *Climate Dyn.*, **16**, <https://doi.org/10.1007/s00382-022-06640-3>
- 849 Miettinen, A., D. Divine, N. Koc, F. Godtlielsen, and I. R. Hall, 2012: Multicentennial variability of  
850 the sea surface temperature gradient across the subpolar North Atlantic over the last 2.8 kyr. *J.  
851 Climate*, **25**, 4205-4219, <https://doi.org/10.1175/JCLI-D-11-00581.1>
- 852 Mikolajewicz, U., and E. Maier-Reimer, 1990: Internal secular variability in an ocean general

- 853 circulation model. *Climate Dyn.*, **4**, 145-156, <https://doi.org/10.1007/bf00209518>
- 854 Mysak, L. A., T. F. Stocker, and F. Huang, 1993: Century-scale variability in a randomly forced, two-  
855 dimensional thermohaline ocean circulation model. *Climate Dyn.*, **8**, 103-116,  
856 <https://doi.org/10.1007/bf00208091>
- 857 Nakamura, M., P. H. Stone, and J. Marotzke, 1994: Destabilization of the thermohaline circulation by  
858 atmospheric eddy transports. *J. Climate*, **7**, 1870-1882, [https://doi.org/10.1175/1520-  
859 0442\(1994\)007<1870:DOTTCB>2.0.CO;2](https://doi.org/10.1175/1520-0442(1994)007<1870:DOTTCB>2.0.CO;2)
- 860 Neale, R. B., and Coauthors, 2010: Description of the Community Atmosphere Model (CAM 4.0).  
861 Technical Report NCAR/TN-485+STR, National Center for Atmospheric Research, Boulder,  
862 CO, 268
- 863 Newby, P. E., B. N. Shuman, J. P. Donnelly, K. B. Karnauskas, and J. Marsicek, 2014: Centennial-to-  
864 millennial hydrologic trends and variability along the North Atlantic coast, USA, during the  
865 Holocene. *Geophys. Res. Lett.*, **41**, 4300-4307, <https://doi.org/10.1002/2014GL060183>
- 866 Nyberg, J., B. A. Malmgren, A. Kuijpers, and A. Winter, 2002: A centennial-scale variability of  
867 tropical North Atlantic surface hydrography during the late Holocene. *Palaeogeogr. Palaeoecol.*  
868 *Palaeoecol.*, **183**, 25-41, [https://doi.org/10.1016/S0031-0182\(01\)00446-1](https://doi.org/10.1016/S0031-0182(01)00446-1)
- 869 Oppo, D. W., J. F. McManus, and J. L. Cullen, 2003: Deepwater variability in the Holocene epoch.  
870 *Nature*, **422**, 277, <https://doi.org/10.1038/422277b>
- 871 Park, W., and M. Latif, 2008: Multidecadal and multicentennial variability of the meridional  
872 overturning circulation. *Geophys. Res. Lett.*, **35**, <https://doi.org/10.1029/2008GL035779>
- 873 Rivin, I., and E. Tziperman, 1997: Linear versus self-sustained interdecadal thermohaline variability  
874 in a coupled box model. *J. Phys. Oceanogr.*, **27**, 1216-1232, [https://doi.org/10.1175/1520-  
875 0485\(1997\)027<1216:LVSSIT>2.0.CO;2](https://doi.org/10.1175/1520-0485(1997)027<1216:LVSSIT>2.0.CO;2)
- 876 Roquet, F., G. Madec, T. J. McDougall, and P. M. Barker, 2015: Accurate polynomial expressions for  
877 the density and specific volume of seawater using the TEOS-10 standard. *Ocean Modell.*, **90**,  
878 29-43, <https://doi.org/10.1016/j.ocemod.2015.04.002>
- 879 Sévellec, F., T. Huck, and M. Ben Jelloul, 2006: On the mechanism of centennial thermohaline  
880 oscillations. *J. Mar. Res.*, **64**, 355-392, <https://doi.org/10.1357/002224006778189608>
- 881 Smith, R. D., and Coauthors, 2010: The Parallel Ocean Program (POP) reference manual. Los Alamos  
882 National Laboratory Tech. Rep. LAUR-10-01853, 140 pp.,  
883 <https://www.cesm.ucar.edu/models/cesm1.0/pop2/doc/sci/POPRefManual.pdf>.
- 884 Stommel, H., 1961: Thermohaline convection with two stable regimes of flow. *Tellus*, **13**, 224-230,

- 885 <https://doi.org/10.1111/j.2153-3490.1961.tb00079.x>
- 886 Suarez, M. J., and P. S. Schopf, 1988: A delayed action oscillator for ENSO. *J. Atmos. Sci.*, **45**, 3283-  
887 3287, [https://doi.org/10.1175/1520-0469\(1988\)045<3283:ADAOFE>2.0.CO;2](https://doi.org/10.1175/1520-0469(1988)045<3283:ADAOFE>2.0.CO;2)
- 888 Sun, D. Z., 1997: El Niño: A coupled response to radiative heating? *Geophys. Res. Lett.*, **24**, 2031-  
889 2034, <https://doi.org/10.1029/97GL01960>
- 890 Sutton, R. T., G. D. McCarthy, J. Robson, B. Sinha, A. T. Archibald, and L. J. Gray, 2018: Atlantic  
891 multidecadal variability and the U.K. ACSIS program. *Bull. Amer. Meteor. Soc.*, **99**, 415-425,  
892 <https://doi.org/10.1175/BAMS-D-16-0266.1>
- 893 Thirumalai, K., T. M. Quinn, Y. Okumura, J. N. Richey, J. W. Partin, R. Z. Poore, and E. Moreno-  
894 Chamarro, 2018: Pronounced centennial-scale Atlantic Ocean climate variability correlated with  
895 Western Hemisphere hydroclimate. *Nat. Commun.*, **9**, 392, [https://doi.org/10.1038/s41467-018-  
896 02846-4](https://doi.org/10.1038/s41467-018-02846-4)
- 897 Thornalley, D. J. R., and Coauthors, 2013: Long-term variations in Iceland-Scotland overflow  
898 strength during the Holocene. *Climate Past*, **9**, 2073-2084, [https://doi.org/10.5194/cp-9-2073-  
899 2013](https://doi.org/10.5194/cp-9-2073-2013)
- 900 Tziperman, E., L. Stone, M. A. Cane, and H. Jarosh, 1994: El Niño chaos: Overlapping of resonances  
901 between the seasonal cycle and the Pacific ocean-atmosphere oscillator. *Science*, **264**, 72-74,  
902 <https://doi.org/10.1126/science.264.5155.72>
- 903 Vellinga, M., and P. L. Wu, 2004: Low-latitude freshwater influence on centennial variability of the  
904 Atlantic thermohaline circulation. *J. Climate*, **17**, 4498-4511, <https://doi.org/10.1175/3219.1>
- 905 Wanner, H., and Coauthors, 2008: Mid- to late Holocene climate change: An overview. *Quat. Sci.*  
906 *Rev.*, **27**, 1791-1828, <https://doi.org/10.1016/j.quascirev.2008.06.013>
- 907 Wei, X., and R. Zhang, 2022: A simple conceptual model for the self-sustained multidecadal AMOC  
908 variability. *Geophys. Res. Lett.*, **49**, 11, <https://doi.org/10.1029/2022GL099800>
- 909 Welander, P., 1982: A simple heat-salt oscillator. *Dyn. Atmos. Oceans*, **6**, 233-242,  
910 [https://doi.org/10.1016/0377-0265\(82\)90030-6](https://doi.org/10.1016/0377-0265(82)90030-6)
- 911 Wills, R. C., T. Schneider, J. M. Wallace, D. S. Battisti, and D. L. Hartmann, 2018: Disentangling  
912 global warming, multidecadal variability, and El Niño in Pacific temperatures. *Geophys. Res.*  
913 *Lett.*, **45**, 2487-2496, <https://doi.org/10.1002/2017GL076327>
- 914 Winton, M., and E. S. Sarachik, 1993: Thermohaline oscillations induced by strong steady salinity  
915 forcing of ocean general-circulation models. *J. Phys. Oceanogr.*, **23**, 1389-1410,  
916 [https://doi.org/10.1175/1520-0485\(1993\)023<1389:TOIBSS>2.0.CO;2](https://doi.org/10.1175/1520-0485(1993)023<1389:TOIBSS>2.0.CO;2)

917 Yin, F. L., 1995: A mechanistic model of ocean interdecadal thermohaline oscillations. *J. Phys.*  
918 *Oceanogr.*, **25**, 3239-3246, [https://doi.org/10.1175/1520-](https://doi.org/10.1175/1520-0485(1995)025<3239:AMMOOI>2.0.CO;2)  
919 [0485\(1995\)025<3239:AMMOOI>2.0.CO;2](https://doi.org/10.1175/1520-0485(1995)025<3239:AMMOOI>2.0.CO;2)

920 Zhang, R., and Coauthors, 2019: A review of the role of the Atlantic meridional overturning  
921 circulation in Atlantic Multidecadal Variability and associated climate impacts. *Rev. Geophys.*,  
922 **57**, 316-375, <https://doi.org/10.1029/2019RG000644>

923